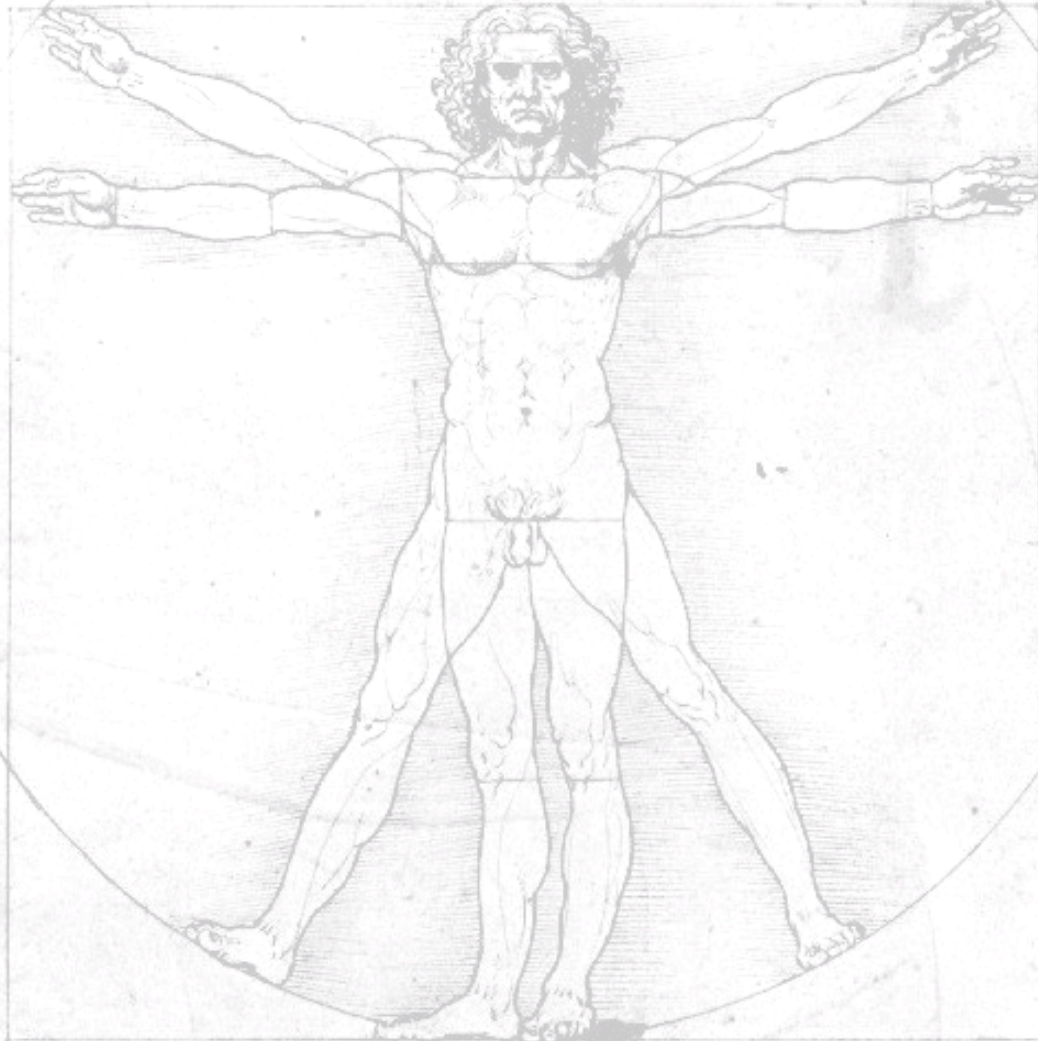


Topological Quantum Computation

Jiannis K. Pachos

Properties of anyons



Bertinoro, June 2013

EPSRC

Engineering and Physical Sciences
Research Council



UNIVERSITY OF LEEDS

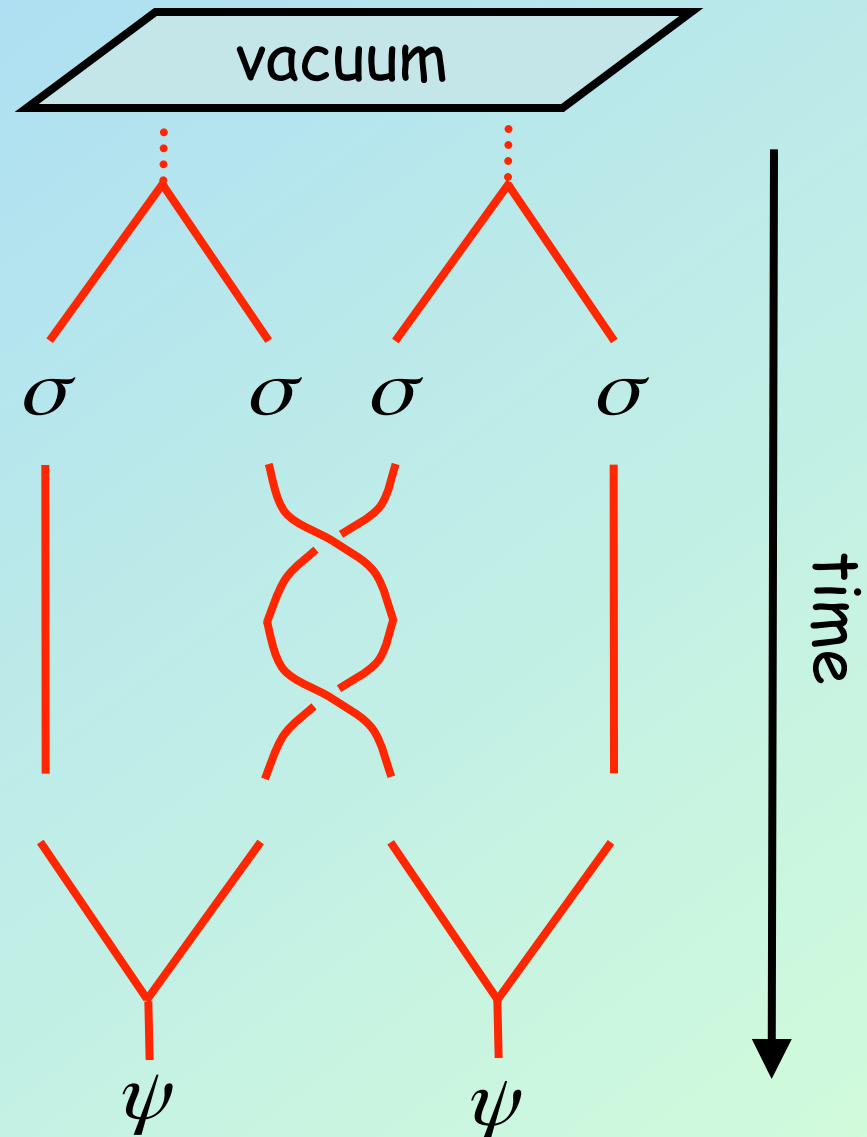
Use anyon for QC

- Assume we can:
 - **Create** identifiable anyons
vacuum pair creation
 - **Braid** anyons
Statistical evolution:
braid representation B
 - **Fuse** anyons

e.g. $\sigma \times \sigma = 1 + \psi$

Fusion Hilbert space:

$$|\sigma, \sigma \rightarrow 1\rangle, |\sigma, \sigma \rightarrow \psi\rangle$$



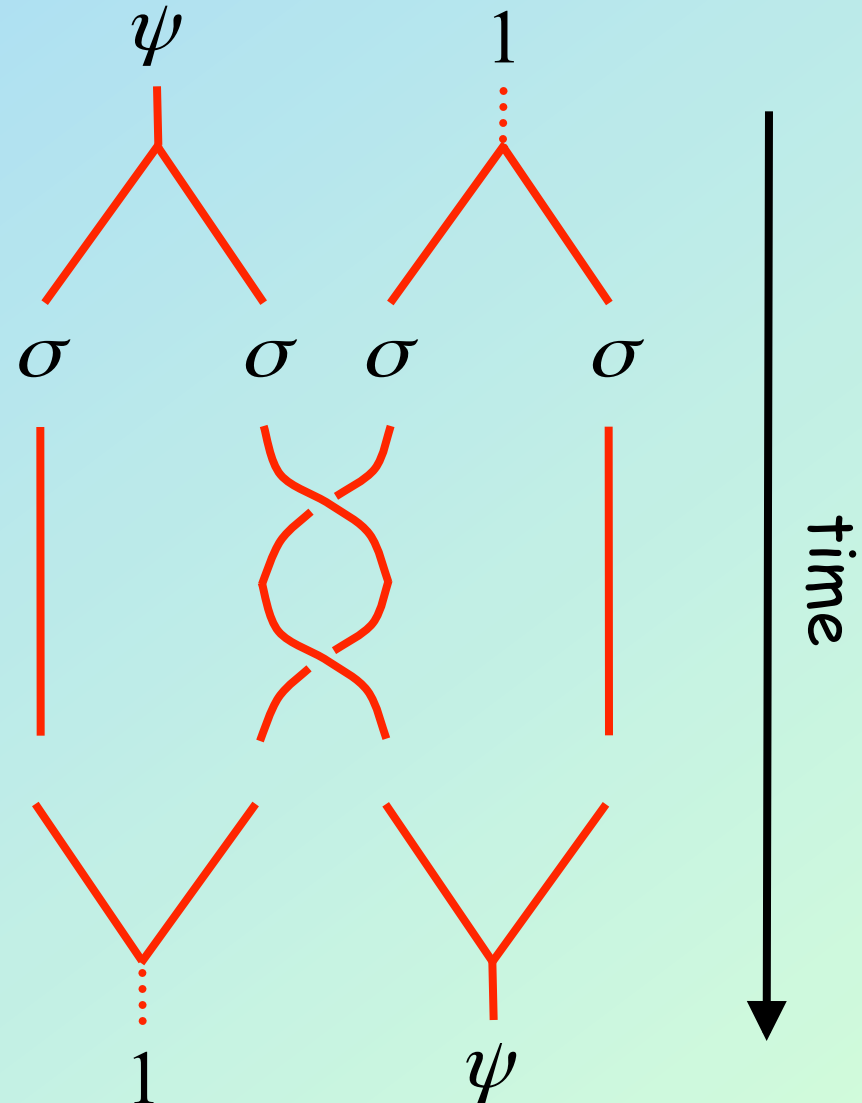
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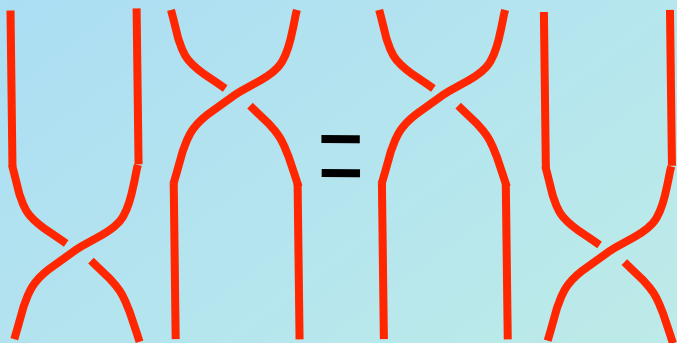
The braid group B_n

The braid group B_n has elements b_1, b_2, \dots, b_{n-1} that satisfy:

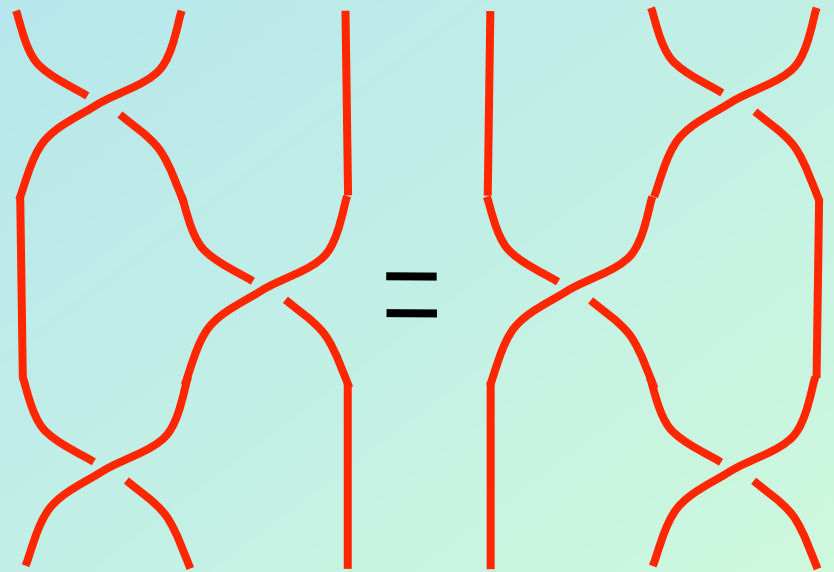
$$b_i b_j = b_j b_i, \text{ for } |i - j| \geq 2$$

$$b_i b_{i+1} b_i = b_{i+1} b_i b_{i+1} \text{ for } 1 \leq i < n$$

Pictorially:



$$b_i b_j = b_j b_i$$

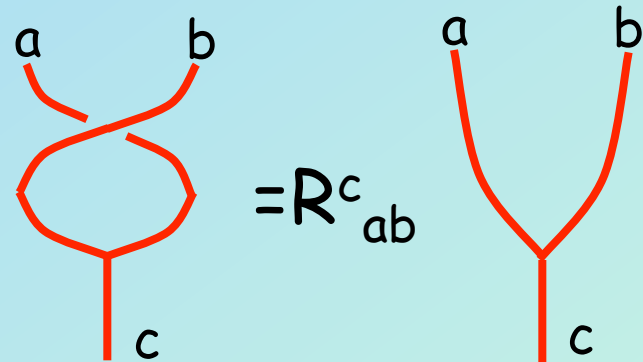


$$b_i b_{i+1} b_i = b_{i+1} b_i b_{i+1}$$

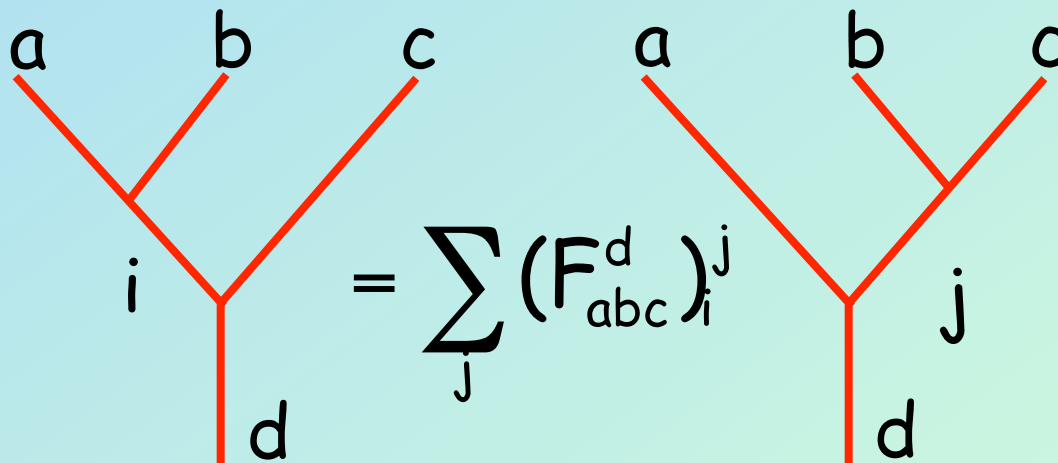
Braiding and Fusion properties

- The action of braiding of two anyons depends on their fusion outcome:

R^c_{ab} is a phase factor



- Changing the order of fusion is non-trivial:



Inception of Anyonic Models

1. Take a certain number of **different anyons**

$1, a, b, \dots$

the vacuum (1) and one or more non-trivial particles

2. Define **fusion rules** between them

$1 \times a = a, a \times b = c + d + \dots, a \times a = 1 + \dots$

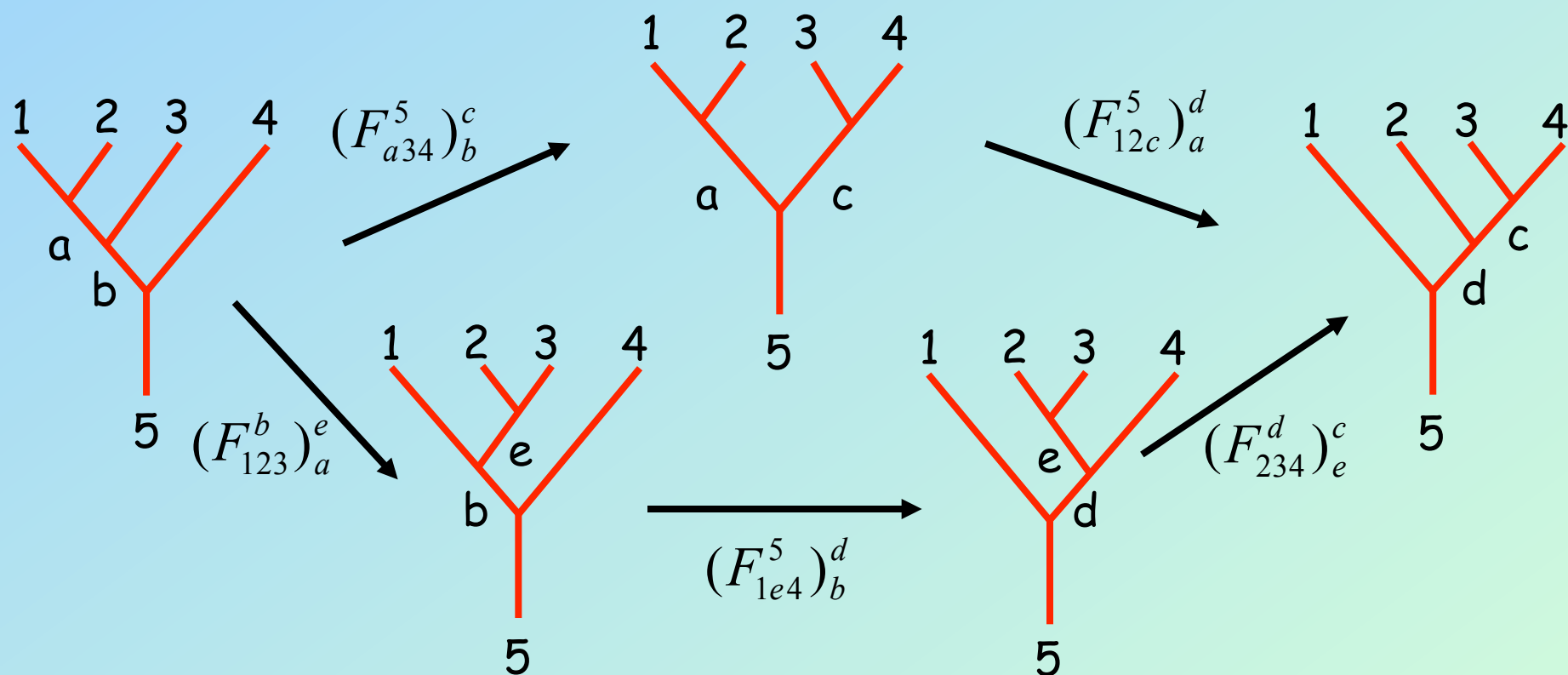
The **vacuum** acts trivially. Each particle has an **anti-particle** (might be itself or not).

- Abelian anyons $a \times b = c$

- Non-Abelian anyons $a \times b = c + d + \dots$

Inception of Anyonic Models

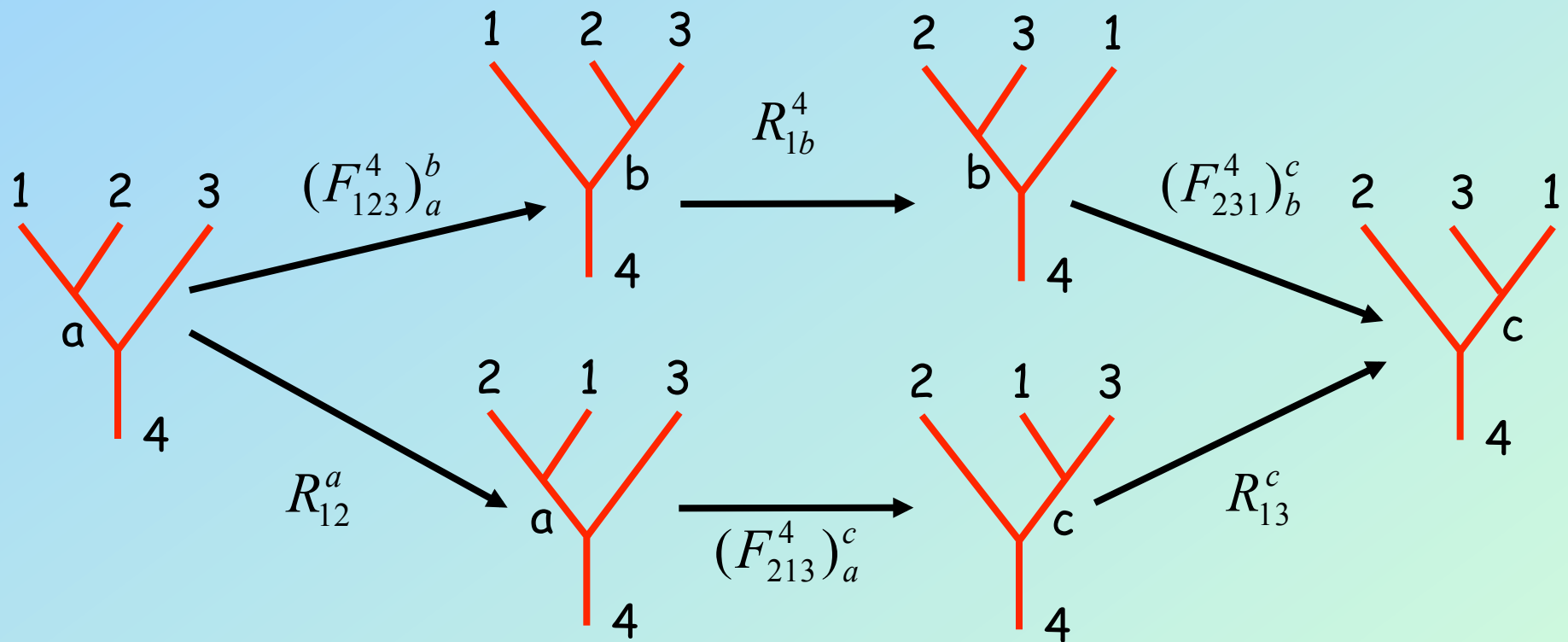
3. The F and B matrices are determined from the **Pentagon** and Hexagon identities



$$(F_{12c}^5)_a^d (F_{a34}^5)_b^c = \sum_e (F_{234}^d)_e^c (F_{1e4}^5)_b^d (F_{123}^b)_a^e$$

Inception of Anyonic Models

3. The F and B matrices are determined from the Pentagon and **Hexagon** identities

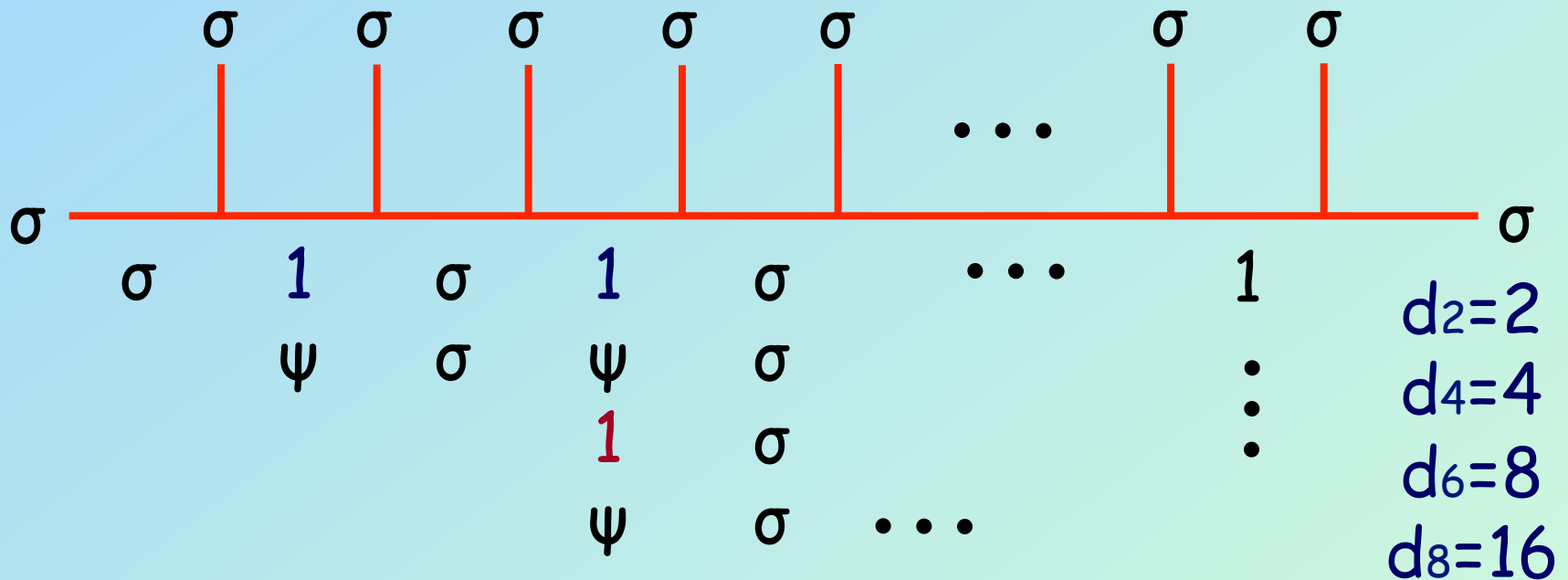


$$\sum_b (F_{231}^4)^c R_{1b}^4 (F_{123}^4)^b = R_{13}^c (F_{213}^4)^c R_{12}^a$$

Ising Anyons

Consider the particles: 1 , σ and ψ

Fusion rules: $\sigma \times \sigma = 1 + \psi$, $\psi \times \psi = 1$, $\sigma \times \psi = \sigma$

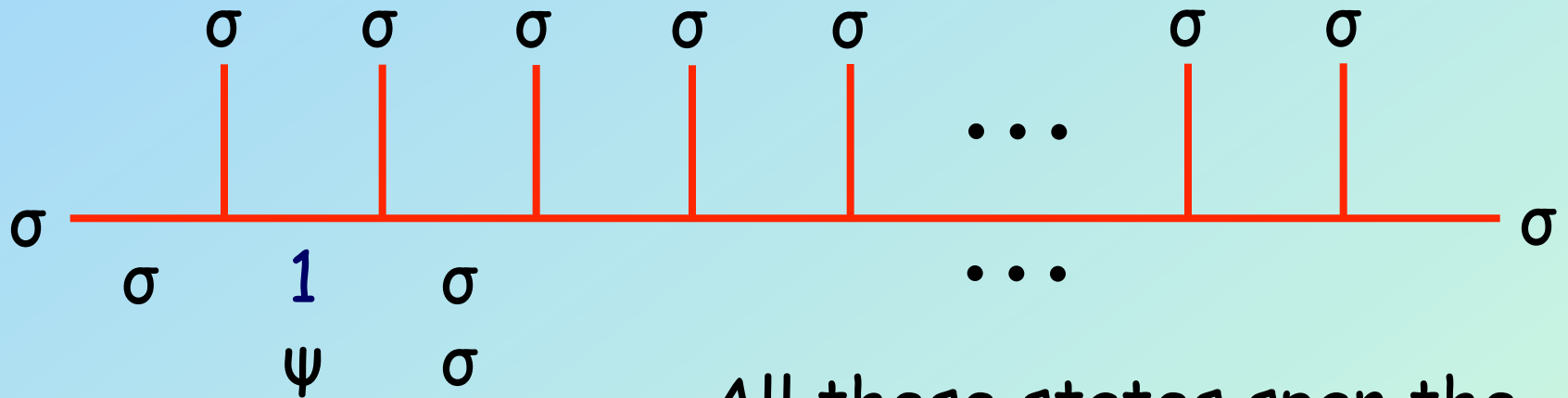


$d_n = 2^{n/2}$ increase in dim of Hilbert space ...

Ising Anyons

Consider the particles: 1 , σ and ψ

Fusion rules: $\sigma \times \sigma = 1 + \psi$, $\psi \times \psi = 1$, $\sigma \times \psi = \sigma$



$$|\Psi\rangle = |1, 1, \dots\rangle$$

$$|\Psi\rangle = |1, \psi, \dots\rangle$$

All these states span the fusion Hilbert space.

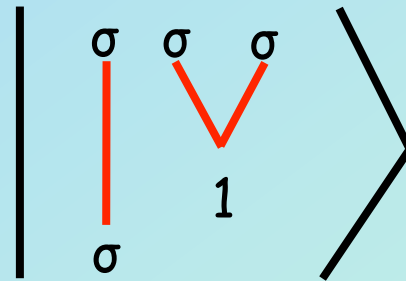
Braiding neighboring anyons transforms states

Ising Anyons

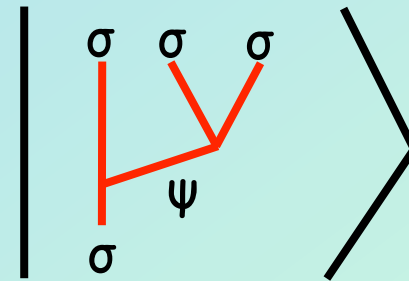
Consider the particles: 1 , σ and ψ

Fusion rules: $\sigma \times \sigma = 1 + \psi$, $\psi \times \psi = 1$, $\sigma \times \psi = \sigma$

Qubit:



State $|0\rangle$



State $|1\rangle$

Ising Anyons

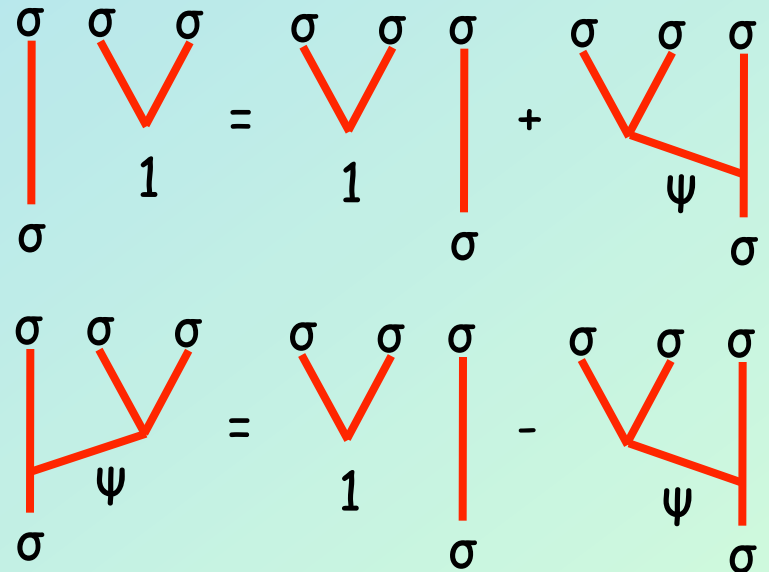
Consider the particles: 1, σ and ψ

Fusion rules: $\sigma \times \sigma = 1 + \psi$, $\psi \times \psi = 1$, $\sigma \times \psi = \sigma$

From 5-gon and 6-gon identities we have:

$$F_{\sigma\sigma\sigma}^{\sigma} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = H$$

Rotation of basis states



Ising Anyons

Braiding $R_{\sigma\sigma}^1 = e^{-i\pi/8}$ and $R_{\sigma\sigma}^\psi = ie^{-i\pi/8} \Rightarrow R_{\sigma\sigma} = e^{-i\pi/8} \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}$

$$(R_{\sigma_1\sigma_2})^2 \begin{array}{c} \sigma_1 \quad \sigma_2 \quad \sigma_3 \\ \text{---} \quad \text{---} \quad \text{---} \\ \text{---} \quad \text{---} \quad \text{---} \\ \sigma_4 \quad \quad \quad 1 \end{array} = (R_{\sigma_1\sigma_2}^1)^2 \begin{array}{c} \sigma_1 \quad \sigma_2 \quad \sigma_3 \\ \text{---} \quad \text{---} \quad \text{---} \\ \text{---} \quad \text{---} \quad \text{---} \\ \sigma_4 \quad \quad \quad 1 \end{array} + (R_{\sigma_1\sigma_2}^\psi)^2 \begin{array}{c} \sigma_1 \quad \sigma_2 \quad \sigma_3 \\ \text{---} \quad \text{---} \quad \text{---} \\ \text{---} \quad \text{---} \quad \text{---} \\ \sigma_4 \quad \quad \quad \psi \end{array}$$

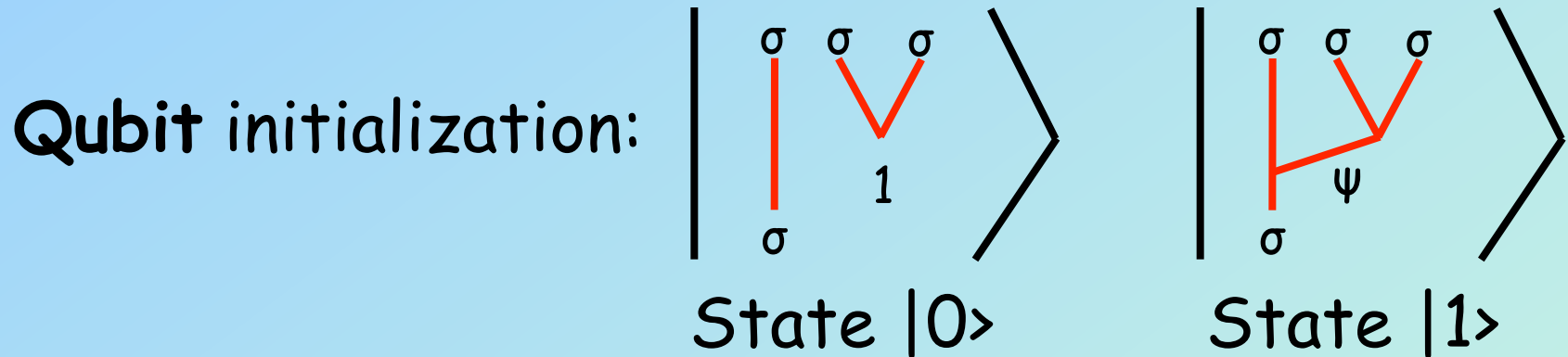
$$= e^{-i\pi/4} \begin{array}{c} \sigma_1 \quad \sigma_2 \quad \sigma_3 \\ \text{---} \quad \text{---} \quad \text{---} \\ \text{---} \quad \text{---} \quad \text{---} \\ \sigma_4 \quad \quad \quad 1 \end{array} - e^{-i\pi/4} \begin{array}{c} \sigma_1 \quad \sigma_2 \quad \sigma_3 \\ \text{---} \quad \text{---} \quad \text{---} \\ \text{---} \quad \text{---} \quad \text{---} \\ \sigma_4 \quad \quad \quad \psi \end{array}$$

$$= e^{-i\pi/4} \begin{array}{c} \sigma_1 \quad \sigma_2 \quad \sigma_3 \\ \text{---} \quad \text{---} \quad \text{---} \\ \text{---} \quad \text{---} \quad \text{---} \\ \sigma_4 \quad \quad \quad \psi \end{array}$$

$$H\sigma^z H = \sigma^x$$

Clifford group:
non-universal!

Ising Anyons



Measurement: Outcome of pairwise fusion, 1 or ψ
 $H\sigma^z H = \sigma^x$

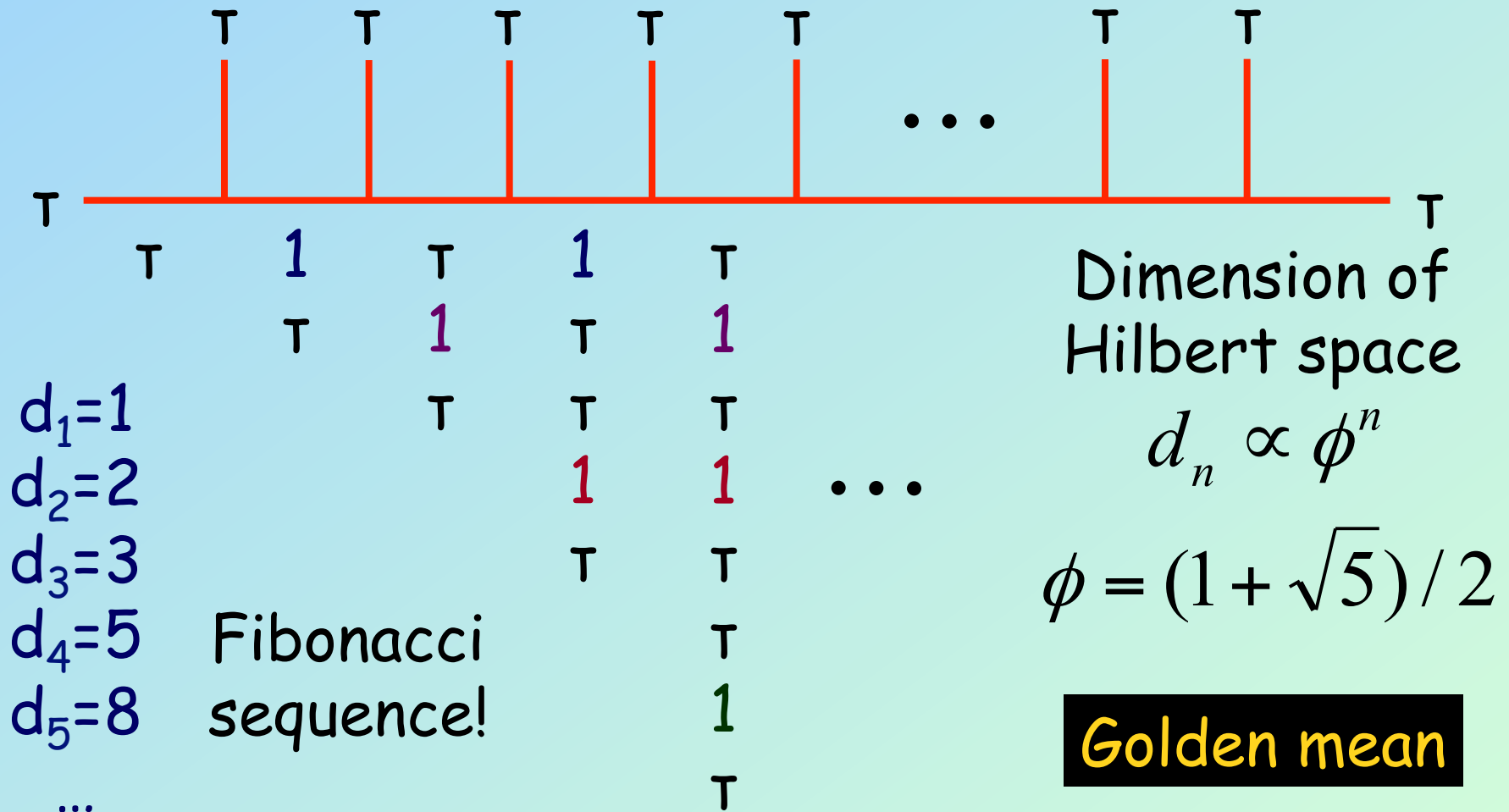
Gates: Clifford group. Non-universal!

One needs a **phase gate**: employ interactions between anyons.

Can be employed as a quantum memory.

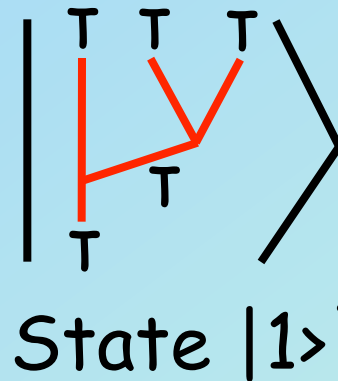
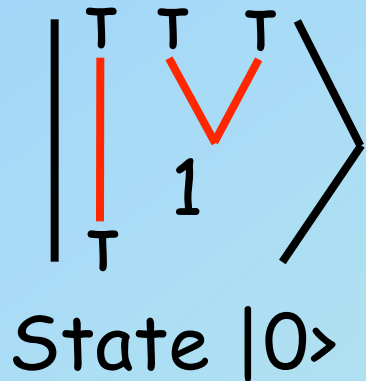
Fibonacci Anyons

Consider anyons with labels 1 or τ with the fusion properties: $1 \times 1 = 1$, $1 \times \tau = \tau$, $\tau \times \tau = 1 + \tau$



Fibonacci Anyons and QC

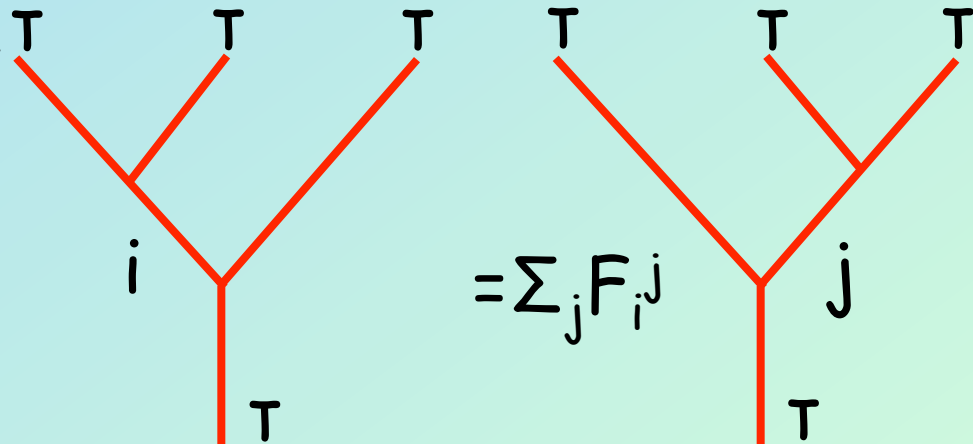
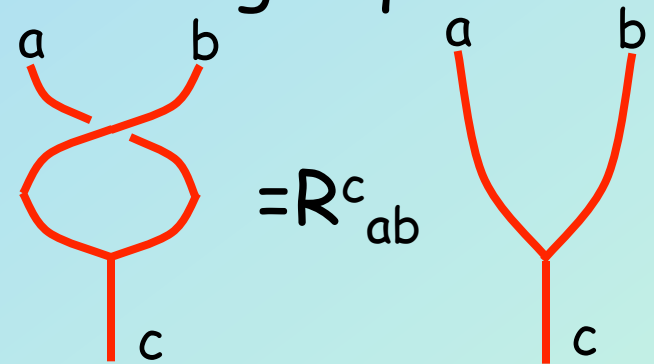
Qubit encoding:



$$|\tau, \tau \rightarrow 1\rangle = |0\rangle$$

$$|\tau, \tau \rightarrow \tau\rangle = |1\rangle$$

Evolving a qubit:

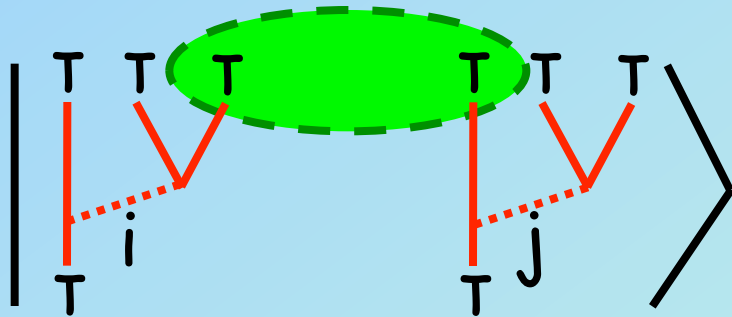


Unitaries B and F are dense in $SU(2)$.

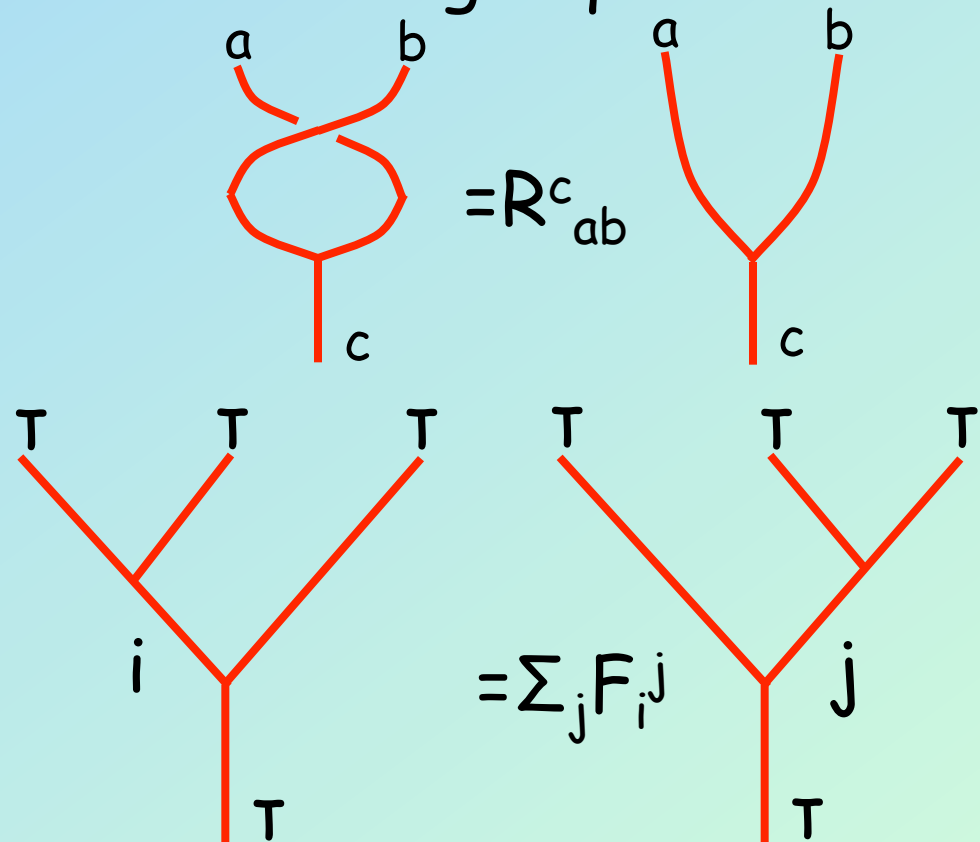
[Freedman, Larsen, Wang, CMP 228, 177 (2002)]

Fibonacci Anyons and QC

Qubit encoding:



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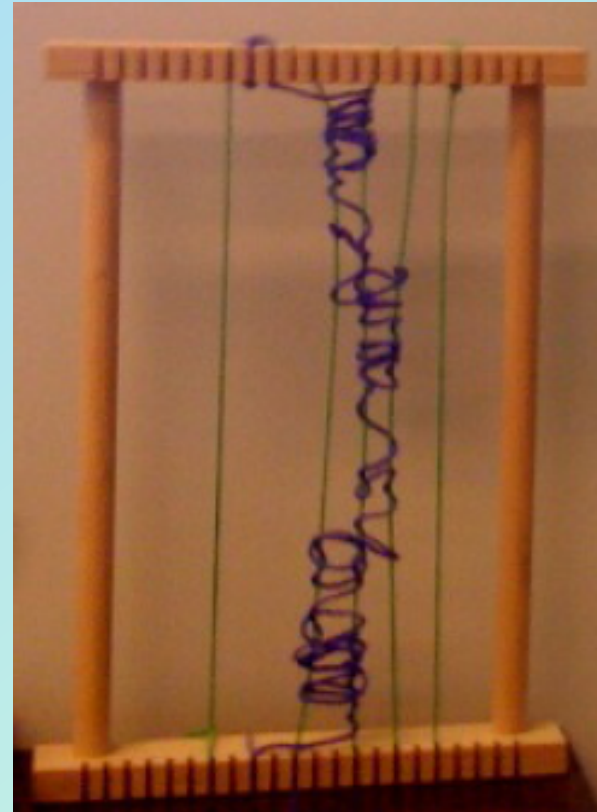
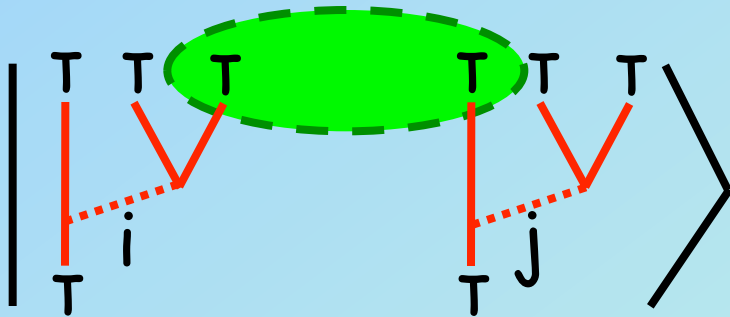


Unitaries B and F are dense in $SU(2)$.

Extends to $SU(d_n)$ when n anyons are employed.

Fibonacci Anyons and QC

Qubit encoding:



CNOT

Unitaries B and F are dense in $SU(2)$.
Extends to $SU(d_n)$ when n anyons are employed.

Conclusions

- Topological Quantum Computation promises to **overcome** the problem of **decoherence** and errors in the most direct way.
- There is lots of work to be done to make anyons work for us.
- Is it worth it?

Aesthetics says YES!

