

# Introduction to Dynamical Systems

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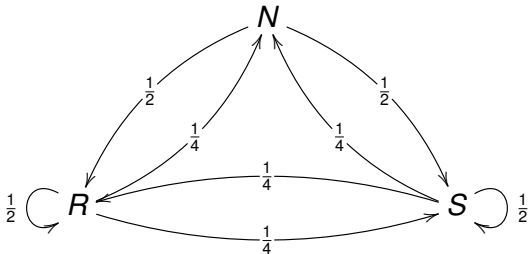
Bertinoro, June 2013

# The Land of Oz

The Land of Oz is blessed with many things, but not by good weather. They never have two nice days in a row. If they have a nice day, the chance of rain or snow the next day are the same. If there is rain or snow the chances are even that the weather stays the same for the next day. If there is a change from snow or rain, only half of the time is this a change to a nice day.

# The Land of Oz

From this we obtain the transition probabilities between nice (N), rainy (R) and snowy (S) days:



# The Land of Oz

We can then define the following **transition matrix**:

$$\mathbf{P} = \begin{pmatrix} \frac{1}{2} & \frac{1}{4} & \frac{1}{4} \\ \frac{1}{2} & 0 & \frac{1}{2} \\ \frac{1}{4} & \frac{1}{4} & \frac{1}{2} \end{pmatrix}$$

From Grinstead & Snell: *Introduction to Probability*, p406;  
available as GNU book on <http://www.dartmouth.edu/~chance>

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# Discrete Time Markov Chain

Given a finite set of states  $S = \{s_1, \dots, s_r\}$ .

A **discrete time Markov chain** (DTMC) on  $S$  is defined via a **stochastic matrix**  $\mathbf{P}$  as above, i.e. an  $r \times r$  (square) matrix with entries  $0 \leq p_{ij} \leq 1$  and such that all row sums are equal to one, i.e.

$$\sum_j p_{ij} = 1.$$

# Discrete Time Markov Processes

Let  $\mathbf{P}$  be the **transition matrix** of a DTMC. The entry in  $p_{ij}^{(n)}$  in the  $n$ -th matrix power  $\mathbf{P}^n$  gives the probability that the Markov chain, starting in state  $s_i$ , will be in state  $s_j$  after exactly  $n$  steps.

At any time step we can describe the probabilities of being in a certain state  $s_i$  by a probability  $u_i$ . These probabilities define a **probability distribution**, i.e. a row vector

$$\mathbf{u} = (u_1, u_2, \dots, u_r)$$

such that  $0 \leq u_i \leq 1$  and  $\sum_i u_i = 1$ .

For any stochastic matrix  $\mathbf{P}$  and probability distribution  $\mathbf{u}$  the multiplication  $\mathbf{uP}$  is again a probability distribution.

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# The Land of Oz

Consider the **initial probability distributions**  $\mathbf{u} = (0, 1, 0)$  and  $\mathbf{v} = (\frac{1}{3}, \frac{1}{3}, \frac{1}{3})$  in the Oz Example. The vector  $\mathbf{u}$  describes a situation where we are certain that we start with a nice day (N), while  $\mathbf{v}$  corresponds to one where we assume the same chances of having a rainy (R), nice (N) or snowy (S) day.

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$$\mathbf{uP} = \left(\frac{1}{2}, 0, \frac{1}{2}\right) \quad \mathbf{uP}^2 = \left(\frac{3}{8}, \frac{1}{4}, \frac{3}{8}\right)$$

# The Land of Oz

Consider the **initial probability distributions**  $\mathbf{u} = (0, 1, 0)$  and  $\mathbf{v} = (\frac{1}{3}, \frac{1}{3}, \frac{1}{3})$  in the Oz Example.

$$\mathbf{vP} = \begin{pmatrix} 0.41667 \\ 0.16667 \\ 0.41667 \end{pmatrix}^T \quad \mathbf{vP}^2 = \begin{pmatrix} 0.39583 \\ 0.20833 \\ 0.39583 \end{pmatrix}^T \quad \mathbf{vP}^3 = \begin{pmatrix} 0.40104 \\ 0.19792 \\ 0.40104 \end{pmatrix}^T$$

$$\mathbf{vP}^4 = \begin{pmatrix} 0.39974 \\ 0.20052 \\ 0.39974 \end{pmatrix}^T \quad \dots \quad \mathbf{vP}^{100} = \begin{pmatrix} 0.40000 \\ 0.20000 \\ 0.40000 \end{pmatrix}^T$$

# Convention

Note that in the theory of Markov chains one usually is concerned with probability distributions as row vectors. Therefore, probability vectors are **post-multiplied** by the stochastic matrix  $\mathbf{P}$  defining a Markov chain.

The usual **pre-multiplication** could be realised via:

$$\mathbf{P}\mathbf{u} = (\mathbf{u}^T \mathbf{P}^T)^T$$

# Dynamical Systems

# Dynamical Systems (Birkhoff 1927)

**Introductory remarks.** In dynamics we deal with physical systems whose state at time  $t$  is completely specified by the values of  $n$  real variables

$$x_1, x_2, \dots, x_n.$$

Accordingly the system is such that the rates of change of these variables, namely

$$dx_1/dt, dx_2/dt, \dots, dx_n/dt,$$

merely depend upon the values of the variables themselves, so that the laws of motion can be expressed by means of  $n$  differential equations of the first order

$$dx_i/dt = X_i(x_1, x_2, \dots, x_n) \quad (i = 1, \dots, n).$$

George D. Birkhoff. *Dynamical Systems*, volume 9 of *Colloquium Publications*. AMS, 1927.

# Dynamical Systems (Bhatia/Szegö 1970)

... the symbol  $X$  denotes a **metric space** [...] and  $R$  stands for the set of real numbers.

**1.1 Definition.** A **dynamical system** on  $X$  is a triplet  $(X, R, \pi)$ , where  $\pi$  is a map from the product space  $X \times R$  into the space  $X$  satisfying the following axioms:

1.1.1  $\pi(x, 0) = x$  for every  $x \in X$  (identity axiom),

1.1.2  $\pi(\pi(x, t_1), t_2) = \pi(x, t_1 + t_2)$  for every  $x \in X$  and  $t_1, t_2 \in R$  (group axiom),

1.1.3  $\pi$  is continuous (continuity axiom).

Nam Parshad Bhatia and Giorgio P. Szegö. *Stability Theory of Dynamical Systems*, volume 161 of *Grundlehren der mathematischen Wissenschaften*.



# Dynamical Systems

## Definition

A general **dynamical system** is a triple  $(G, \pi, X)$  with  $(G, \cdot)$  a group,  $X$  any set and  $\pi : G \times X \rightarrow X$  with:

### Identity Axiom

$$\pi(e, x) = x$$

for all  $x \in X$  and  $e \in G$  unit.

### Homomorphism Axiom

$$\pi(g, \pi(h, x)) = \pi(gh, x)$$

for all  $x \in X$  and  $g, h \in G$ .

# Elements of a General Dynamical System

A general dynamical system is made up of three ingredients:

**Phase Space:** a set  $X$  where “things happen”. This can have additional structure (topology, norm, etc.)

**Phase Group:** the group  $G$  which allows us to “combine” the partial dynamics to obtain a global picture.

**Group Action:** the way in which the dynamics of the group  $G$  is implemented on the phase space  $X$ .

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# Variations of Dynamical System

Typical phase groups are  $\mathbb{Z}$  (integers) or  $\mathbb{R}$  (reals) for so called **discrete time** or **continuous time** models.

To investigate, for example, symmetries of the phase space it is also often the case that one considers so-called **Lie Groups** as transformation groups.

Typically we will request that the group action preserves the structure of the phase space, i.e.  $\pi(g, \cdot)$  is a **structure preserving** morphism on  $X$  for all  $g \in G$ .

An option is to drop invertability to get **one-sided** dynamical systems by taking  $G$  to be a semi-group (e.g. the naturals  $\mathbb{N}$ ).

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# (Phase) Groups

## Definition

A **group**  $G$  is a set with two maps (product and inverse)

$$\cdot : G \times G \rightarrow G \text{ and } \cdot^{-1} : G \rightarrow G$$

fulfilling:

- (i)  $(xy)z = x(yz)$  for all  $x, y, z \in G$  **associativity axiom**.
- (ii)  $\exists e \in G$  such that  $ex = xe = x$  for all  $x \in G$  **identity axiom**.
- (iii)  $x^{-1}x = xx^{-1} = e$  for all  $x \in G$  **inverse axiom**.

Here the group is presented **multiplicatively**, some groups are represented **additively**, e.g.  $(\mathbb{Z}, +)$  and  $(\mathbb{R}, +)$ .

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# G-Spaces

## Definition

Let  $G$  be a group. A **G-Space** is a set  $S$  and a map  $\tau : G \times S \rightarrow S$  so that

$$\tau(e, s) = s \quad \text{all } s \in S$$

and

$$\tau(g, \tau(h, s)) = \tau(gh, s)$$

for all  $g, h \in G$  and  $s \in S$ .  $\tau$  is also called an **action of  $G$  on  $S$** .

We write  $\tau_g(s) = \tau(g, s)$  so  $\tau_g : S \rightarrow S$  and we have  $\tau_g \tau_h = \tau_{gh}$  as well as  $\tau_g \tau_{g^{-1}} = \tau_{g^{-1}} \tau_g = \tau_e = \text{id}$ , see e.g. [3, I.2]

# Phase Spaces

There are many choices for the **phase space** of a dynamical system, among them we could mention:

**Topological Spaces** and require that  $\pi(g, \cdot)$  a homeomorphism.

**Measurable Spaces** with  $\pi(g, \cdot)$  to be measure preserving.

**Vector Spaces** like  $\mathbb{R}^n$  with  $\pi$  a linear map or operator.

**Strings of Symbols** in an alphabet  $\Sigma$  as in Symbolic Dynamics.

**Differentiable Manifolds** as, e.g., in Classical Mechanics.

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# Group Action

## Definition

Let  $(G, \pi, X)$  be a dynamical system. The **orbit** of a point  $x \in X$  is given by

$$O_G(x) = \{\pi(g, x) \mid g \in G\}.$$

## Definition

Let  $(G, \pi, X)$  be a dynamical system. The group action  $\pi$  is  
transitive iff

$$\forall x, x' \in X: O_G(x) = O_G(x').$$

faithful iff

$$g \mapsto \pi(g, x) \text{ is injective.}$$

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# Elements of Ergodic Theory



# Topological Dynamical System

## Definition

A **topological dynamical system** is a dynamical system  $(G, \pi, X)$  with the elements:

$G$  is a topological group, i.e.  $\cdot$  is continuous,

$X$  is a topological space,

and  $\pi$  fulfills the

**Continuity Axiom:**

$$\pi : G \times X \rightarrow X \text{ is continuous.}$$

# Topological Spaces

## Definition

A **topological space** is a set  $X$  together with a family of sub-sets  $\tau \subseteq \mathcal{P}(X)$ , the **topology** (of open sets), iff

- 1  $\emptyset \in \tau$  and  $X \in \tau$ .
- 2  $\bigcap_{i=0}^n O_i \in \tau$  for  $O_i \in \tau$  (finite).
- 3  $\bigcup_{i \in I} O_i \in \tau$  for  $O_i \in \tau$  (arbitrary).

The sets  $O \in \tau$  are called **open** sets. The complements  $A = X \setminus O$  of open sets are **closed** sets.

# Metric Spaces

## Definition

A **metric space** is a set  $X$  and a real-valued function  $d(., .)$ , a **metric**, on  $X \times X$  which satisfies:

- 1  $d(x, y) \geq 0$
- 2  $d(x, y) = 0 \iff x = y$
- 3  $d(x, y) = d(y, x)$
- 4  $d(x, z) \leq d(x, y) + d(y, z)$

# Complete Metric Spaces

In a metric space we can define a basis for the topology open sets via **open balls**, i.e. sets  $B(x, \varepsilon) = \{x' \mid d(x, x') < \varepsilon\}$ , i.e. open sets are those which are unions of open balls.

Given a **sequence**  $(x_i)_{i \in \mathbb{N}}$  of points in a topological space. We say that it **converges** if there exists  $x = \lim x_i$  such that for all neighbourhoods  $U(x)$  of  $x$  there  $\exists N$  s.t. for  $n > N : x_n \in U(x)$ .

A sequence of elements  $(x_i)_{i \in \mathbb{N}}$  in a metric space  $(X, d)$  is called a **Cauchy sequence** if

$$\forall \varepsilon > 0 \exists N : n, m \geq N \Rightarrow d(x_n, x_m) < \varepsilon.$$

A metric space  $(X, d)$  in which all Cauchy sequences converge is called **complete** (metric) space.

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# Continuous Functions

## Definition

A function  $\mathbf{T} : X \rightarrow X'$  between two topological spaces  $(X, \tau)$  and  $(X', \tau')$  is called

continuous iff

$$\forall O \in \tau' : \mathbf{T}^{-1}(O) \in \tau.$$

homeomorph iff

$\mathbf{T}$  is a bijection, and  $\mathbf{T}$  and  $\mathbf{T}^{-1}$  are continuous.

Continuous functions preserve limits, i.e.  $\lim \mathbf{T}(x_i) = \mathbf{T}(\lim(x_i))$ .



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# Measure Theoretical Dynamical System

## Definition

A **measure theoretic dynamical system** is a dynamical system  $(G, \pi, X)$  with

$G$  is a measurable space,

$X$  is a measurable space,

and  $\pi$  fulfills the

**Measurability Axiom:**

$\pi : G \times X \rightarrow X$  is measurable.

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# Measureable Spaces

## Definition

Given any set  $X$ . A family  $\sigma$  of sub-sets  $\sigma \subseteq \mathcal{P}(X)$  is called a  **$\sigma$ -algebra** iff

- 1  $\emptyset \in \sigma$  and  $X \in \sigma$ .
- 2  $\bigcap_{i=0}^{\infty} S_i \in \sigma$  for  $S_i \in \sigma$  (countable).
- 3  $X \setminus S \in \sigma$  for  $S \in \sigma$ .

We say that  $(X, \sigma)$  is a **measurable space**, and  $S \in \sigma$  are **measurable sets**.

By de Morgan we have also:  $\bigcup_{i=0}^{\infty} S_i \in \sigma$  for  $S_i \in \sigma$  (countable).

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# Measures and Measurable Functions

## Definition

Given a measurable space  $(X, \sigma)$  then  $\mu : \sigma \rightarrow \mathbb{R}^+$  is a (finite) **measure** if

①  $\mu(\emptyset) = 0$  (for  $\mu(X) = 1$  we have a **probability measure**).

②  $\mu\left(\bigcup_{i=0}^{\infty} S_i\right) = \sum_{i=0}^{\infty} \mu(S_i)$  for  $S_i \in \sigma$  with  $S_i \cap S_j = \emptyset$  for  $i \neq j$ .

## Definition

A function  $\mathbf{T} : X \rightarrow X'$  between two measure spaces  $(X, \sigma, \mu)$  and  $(X', \tau', \mu')$  is called **measurable** iff

$$\forall S \in \sigma' : \mathbf{T}^{-1}(S) \in \sigma.$$

**measure preserving** iff  $\forall S \in \sigma'$  also  $\mu'(S') = \mu(\mathbf{T}^{-1}(S))$ .



# Topological Mixing Notions

## Definition

Given a topological dynamical system  $(G, \pi, X)$ .

We say that  $(G, \pi, X)$  is **topologically transitive** if

$$\exists x \in X : O_G(x) \text{ is dense in } X.$$

We say that  $(G, \pi, X)$  is (topologically) **minimal** if

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## Definition

A discrete topological dynamical system  $(\mathbf{T}, X)$  is called

**topologically (strong) mixing** if

$$\forall U, V \subseteq X \text{ open and non-empty } \exists N : \forall n > N : \mathbf{T}^n(U) \cap V \neq \emptyset.$$

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# Topologically Transitive

## Theorem

*Given a discrete topological dynamical system  $(\mathbf{T}, X)$  on a compact metric space  $X$  then the following conditions are equivalent:*

- ①  $\forall x \in X : O_{\mathbf{T}}(x)$  is dense in  $X$  (topologically transitive).
- ②  $\forall C \subseteq X$  closed with  $\mathbf{T}(C) = C \Rightarrow C = X$  or  $C = \emptyset$ .
- ③  $\forall O \subseteq X$  open with  $\mathbf{T}(O) = O \Rightarrow O = X$  or  $O = \emptyset$ .
- ④  $\forall O \subseteq X$  open and non-empty, then  $\bigcup_{n=-\infty}^{\infty} \mathbf{T}^n(O) = X$ .

# Measure Theoretic Mixing Notions

## Definition

Given a discrete measure theoretic dynamical system  $(\mathbf{T}, X)$ . We say  $(\mathbf{T}, X)$  is **measure theoretic transitive** or **ergodic** if

$$\forall S \subseteq X \text{ measurable with } \mathbf{T}(S) = S \Rightarrow \mu(S) = 0 \text{ or } \mu(S) = 1.$$

## Definition

A discrete measure theoretic dynamical system  $(\mathbf{T}, X)$  is called **strong mixing** if

$$\forall S_1, S_2 \subseteq X \text{ measurable } \lim_{n \rightarrow \infty} \mu(\mathbf{T}^{-n}(S_1) \cap S_2) = \mu(S_1)\mu(S_2).$$

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## Theorem

*Given a discrete measure theoretic dynamical system  $(\mathbf{T}, X)$  with  $\mathbf{T}$  measure preserving. Then the following conditions are equivalent (with ergodic):*

- ①  $\forall S \subseteq X$  measurable  $\mathbf{T}(S) = S \Rightarrow \mu(S) = 0$  or  $\mu(S) = 1$ .
- ②  $\forall S \subseteq X$  measurable and  $\mu(\mathbf{T}^{-1}(S) \triangle S) = 0$   
 $\Rightarrow \mu(S) = 0$  or  $\mu(S) = 1$ .
- ③  $\forall S \subseteq X$  measurable and  $\mu(S) > 0 \Rightarrow \mu\left(\bigcup_{n=-\infty}^{\infty} \mathbf{T}^{-n}(S)\right) = 1$ .
- ④  $\forall S_1, S_2 \subseteq X$  measurable and  $\mu(S_1) > 0 < \mu(S_2)$   
 $\Rightarrow \exists n \in \mathbb{N}$  such that  $\mu(\mathbf{T}^{-n}(S_1) \cap S_2) = 0$ .

# Ergodic Theorem

Given a discrete measure theoretic dynamical system  $(\mathbf{T}, X)$  and a function (i.e. a random variable)  $f : X \rightarrow \mathbb{R}$ .

The **phase average** of  $f$  is defined as  $\mu(f) = \int_X f(x) dx$ . The **time average** of  $f$  is defined as  $f^*(x) = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T f(\mathbf{T}^t(x)) dt$ .

## Theorem (Birkhoff)

*Given a discrete measure theoretic dynamical system  $(\mathbf{T}, X)$ , with  $\mathbf{T}$  measure preserving, and a function  $f : X \rightarrow \mathbb{R}$  with  $f \in L^1(X, \mu)$  then the following holds:*

$(\mathbf{T}, X)$  is ergodic  $\Leftrightarrow \mu(f) = f^*(x)$   $\mu$ -almost everywhere.

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# Elements of Linear Dynamical Systems

# Linear Dynamical System

## Definition

A **linear dynamical system** is a dynamical system  $(G, \pi, X)$  with

$G$  is a group (typically  $G = \mathbb{Z}$ ),

$X$  is a vector space

and  $\pi$  fulfils the

**Linearity Axiom:**

$$\pi(g, \cdot) : X \rightarrow X \text{ is linear } \forall g \in G.$$

Many versions of linear dynamical systems play an important role in **control theory** investigating e.g. feed back loops etc.

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# Abstract Vector Spaces

## Definition

A **Vector Space** (over a field  $\mathbb{K}$ , e.g.  $\mathbb{R}$  or  $\mathbb{C}$ ) is a set  $\mathcal{V}$  together with two operations:

**Scalar Multiplication**  $\cdot : \mathbb{K} \times \mathcal{V} \mapsto \mathcal{V}$

**Vector Addition**  $+. : \mathcal{V} \times \mathcal{V} \mapsto \mathcal{V}$

such that ( $\forall x, y, z \in \mathcal{V}$  and  $\alpha, \beta \in \mathbb{K}$ ):

- |                               |  |
|-------------------------------|--|
| 1 $x + (y + z) = (x + y) + z$ | 1 $\alpha(x + y) = \alpha x + \alpha y$    |
| 2 $x + y = y + x$             | 2 $(\alpha + \beta)x = \alpha x + \beta x$ |
| 3 $\exists 0 : x + 0 = x$     | 3 $(\alpha\beta)x = \alpha(\beta x)$       |
| 4 $\exists -x : x + (-x) = 0$ | 4 $1x = x$ ( $1 \in \mathbb{K}$ )          |

# Tuple Spaces

## Theorem

*All finite dimensional vector spaces are isomorphic to the (finite) Cartesian product of the underlying field  $\mathbb{K}^n$  (i.e.  $\mathbb{R}^n$  or  $\mathbb{C}^m$ ).*

Finite dimensional vectors can always be represented via their coordinates with respect to a given base, e.g.

$$x = (x_1, x_2, x_3, \dots, x_n)$$

$$y = (y_1, y_2, y_3, \dots, y_n)$$

## Algebraic Structure

$$\alpha x = (\alpha x_1, \alpha x_2, \alpha x_3, \dots, \alpha x_n)$$

$$x + y = (x_1 + y_1, x_2 + y_2, x_3 + y_3, \dots, x_n + y_n)$$

# Linear Operators

## Definition

A map  $\mathbf{T} : \mathcal{V} \rightarrow \mathcal{W}$  between two vector spaces  $\mathcal{V}$  and  $\mathcal{W}$  is called a **linear** map iff

- ①  $\mathbf{T}(x + y) = \mathbf{T}(x) + \mathbf{T}(y)$  and
- ②  $\mathbf{T}(\alpha x) = \alpha \mathbf{T}(x)$

for all  $x, y \in \mathcal{V}$  and all  $\alpha \in \mathbb{K}$  (e.g.  $\mathbb{K} = \mathbb{C}$  or  $\mathbb{R}$ ).

The set of all linear maps between  $\mathcal{V}$  and  $\mathcal{W}$  is denoted  $\mathcal{L}(\mathcal{V}, \mathcal{W})$ . For  $\mathcal{V} = \mathcal{W}$  we talk about a linear **operator** on  $\mathcal{V}$ .

On normed vector spaces the continuous or equivalently **bounded** linear operators are of particular interest, i.e.

$$\mathcal{B}(\mathcal{V}) = \{\mathbf{T} \mid \|\mathbf{T}\| = \sup_{x \in \mathcal{V}} \frac{\|\mathbf{T}(x)\|}{\|x\|} < \infty\} \subseteq \mathcal{L}(\mathcal{V}) = \mathcal{L}(\mathcal{V}, \mathcal{V}).$$

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# Normed Spaces

## Definition

A complex vector space  $\mathcal{V}$  is called a **normed (vector) space** if there is a real valued function  $\|\cdot\|$  on  $\mathcal{V}$  that satisfies ( $\forall x, y \in \mathcal{V}$  and  $\forall \alpha \in \mathbb{C}$ ):

- 1  $\|x\| \geq 0$
- 2  $\|x\| = 0 \iff x = o$
- 3  $\|\alpha x\| = |\alpha| \|x\|$
- 4  $\|x + y\| \leq \|x\| + \|y\|$

The function  $\|\cdot\|$  is called a **norm** on  $\mathcal{V}$ .

We have a **Banach space** if the topology induced by  $d(x, y) = \|x - y\|$  is complete – always for finite dimensional spaces.

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# Hilbert Spaces

## Definition

A complex vector space  $\mathcal{H}$  is called an **inner product space** (or **(pre-)Hilbert space**) if there is a complex valued function  $\langle \cdot, \cdot \rangle$  on  $\mathcal{H} \times \mathcal{H}$  that satisfies ( $\forall x, y, z \in \mathcal{H}$  and  $\forall \alpha \in \mathbb{C}$ ):

- 1  $\langle x, x \rangle \geq 0$
- 2  $\langle x, x \rangle = 0 \iff x = 0$
- 3  $\langle \alpha x, y \rangle = \alpha \langle x, y \rangle$
- 4  $\langle x, y + z \rangle = \langle x, y \rangle + \langle x, z \rangle$
- 5  $\langle x, y \rangle = \overline{\langle y, x \rangle}$

The function  $\langle \cdot, \cdot \rangle$  is called an **inner product** on  $\mathcal{H}$ .

If the topology induced by  $\|x\| = \sqrt{\langle x, x \rangle}$  is complete then we have a **Hilbert space** – always for finite dimensional spaces.

# Basis Vectors

A set of vectors  $x_i$  is said to be **linearly independent** iff

$$\lambda_i x_i = \sum \lambda_i x_i = 0 \quad \text{implies that} \quad \forall i : \lambda_i = 0$$

Two vectors in a Hilbert space are **orthogonal** iff  $\langle x, y \rangle = 0$

An **orthonormal** system (base if it generates all  $\mathcal{H}$ ) in a Hilbert space is a set of linearly independent vectors  $\{b_i\}_i$  with:

$$\langle b_i, b_j \rangle = \delta_{ij} = \begin{cases} 1 & \text{iff } i = j \\ 0 & \text{iff } i \neq j \end{cases}$$

## Theorem

*For a Hilbert space there exists an orthonormal basis  $\{b_i\}$ . The representation of each vector is unique:*

$$x = \sum_i x_i b_i = \sum_i \langle x, b_i \rangle b_i$$

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# Dual Spaces

A **linear functional** on a vector space  $\mathcal{V}$  is a map  $f : \mathcal{V} \rightarrow \mathbb{K}$  such that  $f(x + y) = f(x) + f(y)$  and  $f(\alpha x) = \alpha f(x)$  for all  $x, y \in \mathcal{V}, \alpha \in \mathbb{K}$ .

## Theorem (Riesz Representation Theorem)

*Every (bounded) linear functional on a Hilbert space  $\mathcal{H}$  can be represented by a vector in the Hilbert space  $\mathcal{H}$ , such that*

$$f(x) = \langle y_f | x \rangle = f_y(x)$$

The **dual** Hilbert space  $\mathcal{H}^*$  is isomorphic to the original Hilbert space  $\mathcal{H}$ , e.g. for the universal Hilbert space  $\ell_2(\mathbb{N})^* = \ell_2(\mathbb{N})$ .

$$\ell_p(\mathbb{X}) = \left\{ (x_i)_{i \in \mathbb{X}} \mid \left( \sum_{i \in \mathbb{X}} |x_i|^2 \right)^{\frac{1}{p}} \right\}$$



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# Finite-Dimensional Hilbert Spaces

We represent vectors and their **transpose** using coordinates:

$$\vec{x} = \begin{pmatrix} x_1 \\ \vdots \\ x_n \end{pmatrix}, \quad \vec{y} = \begin{pmatrix} y_1 \\ \vdots \\ y_n \end{pmatrix}^T = (y_1, \dots, y_n)$$

The **adjoint** of  $\vec{x} = (x_1, \dots, x_n)$ , with  $\cdot^* = \bar{\cdot}$  denoting **complex conjugate** in  $\mathbb{C}$ ), is given by

$$\vec{x}^\dagger = \vec{x}^* = (x_1^*, \dots, x_n^*)^T$$

The **inner product** is:

$$\langle \vec{y}, \vec{x} \rangle = \sum_i y_i^* x_i = \vec{y}^\dagger \vec{x}$$

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# Differential Equations

# Discrete (Time) Dynamical Systems: Collatz

The **Collatz problem** is a (one-sided) discrete time dynamical system  $(\mathbf{C}, \mathbb{Z})$ , which we can describe by the following transformation:

$$\mathbf{C} : \mathbb{Z} \rightarrow \mathbb{Z}$$

with

$$\mathbf{C}(n) = \begin{cases} n/2 & \text{if } n \text{ is even} \\ 3 \times n + 1 & \text{otherwise.} \end{cases}$$

The unsolved question is:

Does  $\exists m \in \mathbb{N}$  such that  $\mathbf{C}^m(n) = 1$  for all  $n \in \mathbb{N}$ ?

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# Continuous Dynamical Systems

A popular way to specify continuous time dynamical systems is via (ordinary) differential equations, e.g. [Morris W. Hirsch, Stephen Smale, and Robert L. Devaney. \*Differential Equations, Dynamical Systems and An Introduction to Chaos\*. Elsevier, 2004.](#)

The group action is interpreted as **time**  $t \in \mathbb{R}$ .

## Ordinary Differential Equations

$$\begin{array}{rcl} x_1' & = & \frac{dx_1}{dt} = f_1(t, x_1, x_2, \dots, x_n) \\ x_2' & = & \frac{dx_2}{dt} = f_2(t, x_1, x_2, \dots, x_n) \\ \dots & & \dots \\ x_n' & = & \frac{dx_n}{dt} = f_n(t, x_1, x_2, \dots, x_n) \end{array}$$

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# Differential

Given a function  $f : \mathbb{R} \rightarrow \mathbb{R}$ . We say it is **differentiable** at a point  $t \in \mathbb{R}$  if there is a linear map  $Df(t) : \mathbb{R} \rightarrow \mathbb{R}$  which approximates  $f$  at  $t$ . That is,  $\forall \varepsilon > 0$  there is a neighborhood  $U$  of  $t$  such that:

$$\|f(t') - f(t) - Df(t)(t - t')\| < \varepsilon \|t - t'\| \quad \forall t' \in U$$

We also write for the **differential** (quotient)  $Df = \frac{df}{dt}$ .

We also approximate a function  $f : \mathbb{R}^n \rightarrow \mathbb{R}^m$  by a linear map  $Df(x) : \mathbb{R}^n \rightarrow \mathbb{R}^m$  represented by the matrix of **partial derivatives**:

$$(Df)_{ij} = \frac{\partial f_i}{\partial t_j} \quad i = 1, \dots, m, j = 1, \dots, n.$$

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# Euler's Number

What is Euler's **e**? Metafont users know, it is  **$e = 2.7183\dots$**

It is the unique number such that for  $f(t) = e^t$  we have

$$\frac{df}{dt}(t) = f(t)$$

The exponential function is the fixed-point or eigen-function of the differential operator  $\frac{d}{dt}$ . One could show this via the **Taylor expansion** of  $\exp(x) = \sum_{n=0}^{\infty} \frac{x^n}{n!}$  as  $\frac{d}{dt} \frac{t^n}{n!} = \frac{nt^{n-1}}{n(n-1)!} = \frac{t^{n-1}}{(n-1)!}$ .

The simplest differential equation one can think of is perhaps:

$$x'(t) = \frac{dx}{dt}(t) = ax(t)$$

The solution is  **$x(t) = ke^{at} = k \exp(at)$**  for some constant  $k$  (can be determined via an initial/boundary condition, e.g.  $x(0)$ ).

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# Ordinary Linear Differential Equations [1, p129]

Solution to ordinary differential equations via exponentiation.

$$\begin{aligned}x'_1 &= \frac{dx_1}{dt} = a_{11}x_1 + a_{12}x_2 + \dots + a_{1n}x_n \\x'_2 &= \frac{dx_2}{dt} = a_{21}x_1 + a_{22}x_2 + \dots + a_{2n}x_n \\&\dots \\x'_n &= \frac{dx_n}{dt} = a_{n1}x_1 + a_{n2}x_2 + \dots + a_{nn}x_n\end{aligned}$$

## Theorem

Let  $\mathbf{A}$  be an  $n \times n$  matrix. Then the *unique* solution to the initial value problem  $\mathbf{x}' = \mathbf{A}\mathbf{x}$  with  $\mathbf{x}(0) = \mathbf{x}_0$  is given by

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The exponential of a matrix  $\mathbf{A}$  can be computed as:

$$\exp(\mathbf{A}) = \sum_{k=0}^{\infty} \frac{\mathbf{A}^k}{k!}$$

However this anything but an efficient way to compute it.

We can represent matrices e.g. in Jordan normal form:

$\mathbf{A} = \mathbf{D} + \mathbf{N}$  where  $\mathbf{D}$  is a **diagonal** matrix and  $\mathbf{N}$  is an upper diagonal matrix which is **nilpotent**, i.e.  $\exists m$  s.t.  $\mathbf{N}^m$  vanishes. This boils down to finding the eigenvalues of  $\mathbf{A}$  (via SVD).

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# Smooth Functions

We say a function  $f$  is **differentiable** on  $\mathbb{R}$  or  $U \in \mathbb{R}$  if it is differentiable at every point  $t \in \mathbb{R}$  or  $t \in U$ . We then write  $f \in \mathbf{C}^1(\mathbb{R}) = C^1$  or  $f \in \mathbf{C}(U)$ .

We can see  $Df(x)$  itself as a function  $\mathbb{R}^n \rightarrow \mathcal{L}(\mathbb{R}^n, \mathbb{R}^m) = \mathbb{R}^{nm}$ .

As such we can ask if this is itself differentiable. We denote the set of  $p$ -times differentiable maps by  $\mathbf{C}^p$  and by  $\mathbf{C}^\infty$  the set of infinitely differentiable or **smooth** functions.

Note: Differentiation is primarily a **real** number notion. We need to introduce the notion of a **differentiable manifold** as a space which looks like  $\mathbb{R}^m$  locally (with respect to diff. operations).

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# Differentiable Manifolds

## Definition

Let  $M$  be a topological space.

A **chart**  $(V, \Phi)$  is a homeomorphism  $\Phi$  of an open set  $V$  of  $M$  into an open set of  $\mathbb{R}^m$ .

Two charts  $(V_1, \Phi_1)$  and  $(V_2, \Phi_2)$  are said to be **compatible** in case  $V_1 \cap V_2 = \emptyset$  or the restricted maps  $\Phi_1 \circ \Phi_2^{-1}$  and  $\Phi_2 \circ \Phi_1^{-1}$  are in  $C^\infty(\mathbb{R}^m)$ .

A **atlas** is a set of compatible charts that cover all of  $M$ . Two atlases are compatible if all their charts are.

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# Tangents

## Definition

Let  $f, g \in \mathcal{L}(\mathcal{V}, \mathcal{W})$  then we say that  $f$  is **tangent** to  $g$  at  $t$  iff

$$\lim_{t' \rightarrow t} \frac{\|f(t') - g(t')\|}{\|t' - t\|} = 0$$

## Definition

Let  $M$  be a manifold and  $m \in M$ . A **curve** at  $m$  is a  $C^1$  map  $c : I \rightarrow M$  with an open interval in  $\mathbb{R}$  containing 0 s.t.  $c(0) = m$ .

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# Lie Groups

## Definition

A **Lie group** over a field  $\mathbb{K}$  is a group  $G$  equipped with the structure of a differentiable manifold over  $\mathbb{K}$  such that

$$\dots : G \times G \rightarrow G \text{ is differentiable.}$$

Using the implicit function theorem, one can also show that  $g \mapsto g^{-1}$  is differentiable (a diffeomorphism).

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## Definition

A **field** is a set  $\mathbb{K}$  together with two operations:

**Addition**  $.+. : \mathbb{K} \times \mathbb{K} \mapsto \mathbb{K}$

**Multiplication**  $.\cdot. : \mathbb{K} \times \mathbb{K} \mapsto \mathbb{K}$

$$\textcircled{1} \quad \forall x, y, z \in \mathbb{K} : x + (y + z) = (x + y) + z$$

$$\textcircled{2} \quad \exists 0 \in \mathbb{K}, \forall x \in \mathbb{K} : 0 + x = x$$

$$\textcircled{3} \quad \forall x \in \mathbb{K}, \exists -x \in \mathbb{K} : x + (-x) = 0$$

$$\textcircled{4} \quad \forall x, y \in \mathbb{K} : x + y = y + x$$

such that

$$\textcircled{5} \quad \forall x, y, z \in \mathbb{K} : x \cdot (y \cdot z) = (x \cdot y) \cdot z$$

$$\textcircled{6} \quad \exists 0 \neq e \in \mathbb{K}, \forall x \in \mathbb{K} : e \cdot x = x$$

$$\textcircled{7} \quad \forall 0 \neq x \in \mathbb{K}, \exists x^{-1} \in \mathbb{K} : x \cdot x^{-1} = e$$

$$\textcircled{8} \quad \forall x, y \in \mathbb{K} : x \cdot y = y \cdot x$$

$$\textcircled{9} \quad x \cdot (y + z) = x \cdot y + x \cdot z, \quad \forall x, y, z \in \mathbb{K}$$

# Examples of Lie Groups

Examples of Lie groups we can mention here:

- The additive group of the field  $\mathbb{K} = \mathbb{K}^+$ .
- The multiplicative group of the field  $\mathbb{K}^\times$ .
- The “circle”  $\mathbb{T} = \{z \in \mathbb{C} \mid |z| = 1\}$  or  $\{e^{i\phi} \mid \phi \in [0, 2\pi)\}$ .
- $GL_n(\mathbb{K})$  of invertible matrices of order  $n$  over  $\mathbb{K}$ .
- $SL_n(\mathbb{K})$  of matrices of order  $n$  over  $\mathbb{K}$  with  $\det = 1$ .
- $O_n(\mathbb{K})$  orthogonal matrices over  $\mathbb{K}$  of order  $n$ .
- $U_n$  unitary matrices over  $\mathbb{C}$  of order  $n$ .
- $SO_n(\mathbb{K}) = O_n(\mathbb{K}) \cap SL_n(\mathbb{K})$ .
- $SU_n = U_n \cap SL_n(\mathbb{C})$ .

# Lie Algebras

## Definition

A **Lie algebra** is a vector space  $\mathfrak{g}$  over some field  $\mathbb{K}$  together with a binary operation  $[\cdot, \cdot] : \mathfrak{g} \times \mathfrak{g} \rightarrow \mathfrak{g}$  called the **Lie bracket**, which satisfies the following:

**Bilinearity:**  $\forall \alpha, \beta \in \mathbb{K}$  and  $\forall x, y, z \in \mathfrak{g}$

$$[\alpha x + \beta y, z] = \alpha[x, z] + \beta[y, z]$$

$$[z, \alpha x + \beta y] = \alpha[z, x] + \beta[z, y]$$

**Alternating on  $\mathfrak{g}$ :**  $\forall x \in \mathfrak{g}$

$$[x, x] = 0$$

**Jacobi identity:**  $\forall x, y, z \in \mathfrak{g}$

$$[x, [y, z]] + [z, [x, y]] + [y, [z, x]] = 0$$

# Examples of Lie Algebras

It follows easily  $\forall x, y \in \mathfrak{g}$  that  $[x, y] = -[y, x]$ . One could also define an associative product on an algebra  $\mathfrak{g}$  and then introduce the Lie bracket as  $[x, y] = xy - yx$ .

## Theorem

*Given a Lie group  $G$  then the tangent space at the unit  $\mathfrak{g} = T_e G$  is a Lie algebra.*

Let  $g(t)$  and  $h(t)$  be differentiable paths or  $C^1$  curves on  $G$ . Assume,  $g(0) = h(0) = e$  as well as  $\frac{dg}{dt}(0) = \xi$  and  $\frac{dh}{dt}(0) = \eta$  then we define a Lie bracket on the tangent space  $T_e(G)$  via

$$[\xi, \eta] = \frac{\partial^2}{\partial t \partial s} [g(t), h(s)]|_{s=t=0}$$

where  $[g, h] = ghg^{-1}h^{-1}$  is the *group commutator*.

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# Prescribed Velocities

## Definition

A path or curve  $g(t)$  in a Lie group  $G$  with  $t \in \mathbb{R}$  is called a **one-parameter subgroup** if

$$g(t + s) = g(t)g(s).$$

We denote by  $g_\xi(s)$  the one-parameter sub-group with  $g' = \frac{dg}{dt}(s) = \xi(s)$  – i.e. with prescribed “velocity”  $\xi(s)$ .

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For a Lie group  $G$  and  $\xi \in \mathfrak{g}$ , i.e. its Lie algebra, we define:

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# Exponentiation

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*The exponential map  $\exp : \mathfrak{g} \rightarrow G$  maps a neighbourhood of zero in the tangent algebra  $\mathfrak{g} = \mathbf{T}_e(G)$  diffeomorphically onto a neighbourhood of the identity in  $G$ .*

## Theorem

*Let  $\mathfrak{g} = \mathfrak{a}_1 \oplus \dots \oplus \mathfrak{a}_k$  be a decomposition of a Lie algebra as direct sum, then  $\xi_1 + \dots + \xi_k \mapsto \exp(\xi_1) \dots \exp(\xi_k)$  maps a neighbourhood of zero in  $\mathfrak{g}$  diffeomorphically onto a neighbourhood of the identity in  $G$ .*

If  $G$  is the group of invertible elements in an associative algebra (e.g. of non-singular matrices), then

$$\exp(\xi) = \sum_{n=0}^{\infty} \frac{\xi^n}{n!}.$$

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# Stochastic Dynamics



# Discrete Time Markov Chains (DTCM)

## Definition

A **discrete time Markov chain** (DTMC) on  $S$  is defined via a **stochastic matrix**  $\mathbf{P}$ , i.e. an  $r \times r$  (square) matrix with entries  $0 \leq p_{ij} \leq 1$  and such that all row sums are equal to one, i.e.

$$\sum_j p_{ij} = 1.$$

This defines a discrete linear dynamical system:

Phase group:  $\mathbb{Z}$  or  $\mathbb{N}$ ,

Phase space:  $\mathbb{R}^r$ ,

Group action:  $\pi(n, x) = x \cdot \mathbf{P}^n$ .

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# Memoryless Property of DTMC

Let  $I$  be a finite (or maybe countable) set. Each  $i \in I$  is called a state or **index**. Given a probability space  $(\Omega, \sigma, \mathbb{P})$  a **random variable** is a map  $X : \Omega \rightarrow I$ .

A sequence of random variables  $X_n$  is a **Markov Chain** if

$$\begin{aligned}\mathbb{P}(X_{n+1} = i + 1 \mid X_0 = i_0, X_1 = i_1, \dots, X_n = i_n) &= \\ &= \mathbb{P}(X_{n+1} = i + 1 \mid X_n = i_n) = \\ &= p_{i_n, i_{n+1}} \mathbb{P}(X_n = i_n)\end{aligned}$$

The probability  $\mathbb{P}(i \rightarrow^n j)$  of reaching state (actually index)  $j$  from  $i$  in exactly  $n$  steps is given by  $p_{ij}^{(n)}$  i.e. the entry in row  $i$  and column  $j$  of  $\mathbf{P}^n$ .

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# Properties

## Definition

Given a DTMC with transition matrix  $\mathbf{P}$ . A state  $i$  is said to be

**recurrent** if  $\mathbb{P}(i \rightarrow^n i \text{ for infinitely many } n) = 1$

**transient** if  $\mathbb{P}(i \rightarrow^n i \text{ for infinitely many } n) = 0$

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A DTMC with transition matrix  $\mathbf{P}$  is called

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# Long Run Behaviour

## Theorem

Given a DTMC with transition matrix  $\mathbf{P}$ . If it is **regular** and  $v$  an arbitrary probability vector. Then

$$\lim_{n \rightarrow \infty} v\mathbf{P}^n = w$$

where  $w$  is the unique probability vector for  $\mathbf{P}$ .

## Theorem

Given a DTMC with transition matrix  $\mathbf{P}$ . Assume  $\mathbf{P}$  is **ergodic**. Let  $\mathbf{A}_n$  be the matrix defined by:

$$\mathbf{A}_n = \frac{\mathbf{I} + \mathbf{P} + \dots + \mathbf{P}^n}{n + 1}$$

then  $\mathbf{A}_n \rightarrow \mathbf{W}$  where  $\mathbf{W}$  is a matrix all of whose rows are equal to the unique vector  $w$  for  $\mathbf{P}$ .



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# Continuous Time Markov Chains (CTMC)

## Definition

A **continuous time Markov chain** (CTMC) on  $S = \{s_1, \dots, s_r\}$  is defined via an  $r \times r$  (square) **generator** or **Q-matrix**  $\mathbf{Q} = (q_{ij})$  specifying the **rates** going from an index or state  $i$  to an index or state  $j$  and which fulfills:

- 1  $0 \leq -q_{ii} < \infty$  for all  $i$
- 2  $q_{ij} \geq 0$  for all  $i \neq j$
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# Computing Transition Probabilities

Again we use exponentiation to get the transition probabilities.

$$\mathbf{P}(t) = \exp(t\mathbf{Q}) = \sum_{k=0}^{\infty} \frac{(t\mathbf{Q})^k}{k!}$$

This gives the unique solutions to the **forward equations**

$$\frac{d}{dt}\mathbf{P}(t) = \mathbf{P}(t)\mathbf{Q} \text{ with } \mathbf{P}(0) = \mathbf{I}$$

and the **backward equation**

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# Dynamical Systems in Physics

# Classical Mechanics (in 10min)

Consider point particles (no volume) with mass  $m$ . The position of a particle is given in some coordinates  $q_i$ .

The velocity of the particle is given by

$$v_i = \frac{dq_i}{dt} = \dot{q}_i$$

its acceleration is given by

$$a_i = \frac{dv}{dt} = \frac{d^2 q_i}{dt^2} = \ddot{q}_i$$

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# Lagrange Formalism

Describe the dynamics of a mechanical system via the **Lagrange function** or **Lagrangian**

$$L(q_1, q_2, \dots, q_s, \dot{q}_1, \dot{q}_2, \dots, \dot{q}_s, t)$$

the **action** is defined as  $S = \int_{t_1}^{t_2} L(q_i, \dot{q}_i, t) dt$ .

The **Principle of Least Action** then implies the **Lagrange equations** which give the dynamics:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} - \frac{\partial L}{\partial q_i} = 0$$

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Single Particle:

$$L = \frac{m}{2}(\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$$

Pendulum: length  $l$ , angle  $\phi$ , mass  $m$ , gravitational constant  $g$

$$L = \frac{m}{2}l^2\dot{\phi}^2 + mgl\cos(\phi)$$

Double Pendulum: angles  $\phi_1$  and  $\phi_2$ , lengths  $l_1$  and  $l_2$ , masses  $m_1$  and  $m_2$  [2, p13]:

$$\begin{aligned} L = & \frac{m_1 + m_2}{2}l_1^2\dot{\phi}_1^2 + \frac{m_2}{2}l_2^2\dot{\phi}_2^2 + \\ & m_2l_1l_2\dot{\phi}_1\dot{\phi}_2\cos(\phi_1 - \phi_2) + \\ & (m_1 + m_2)gl_1\cos(\phi_1) + m_2gl_2\cos(\phi_2) \end{aligned}$$

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# Hamiltonian Formalism

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$$H(p_i, q_i, t) = \sum_i p_i \dot{q}_i - L(q_i, \dot{q}_i, t)$$

The dynamics of the system is then described via the Hamiltonian or **canonical** equations:

$$\begin{aligned}\dot{q}_i &= \frac{dq}{dt} = \frac{\partial H}{\partial p_i} \\ \dot{p}_i &= \frac{dp}{dt} = -\frac{\partial H}{\partial q_i}\end{aligned}$$

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The dynamics of the system is then described via the Hamiltonian or **canonical** equations:

$$\begin{aligned}\dot{q}_i &= \frac{dq}{dt} = \frac{\partial H}{\partial p_i} \\ \dot{p}_i &= \frac{dp}{dt} = -\frac{\partial H}{\partial q_i}\end{aligned}$$

L.D. Landau and E.M. Lifschitz. *Mechanik*. Akademie-Verlag, Berlin, 1981.

# Hamiltonian Examples

Single Particle:

$$H = \frac{1}{2m}(p_x^2 + p_y^2 + p_z^2)$$

Particle in Field:

$$H = \frac{1}{2m}(p_x^2 + p_y^2 + p_z^2) + U(x, y, z)$$

Pendulum: with  $p_\phi = ml^2\dot{\phi}$  and  $\dot{\phi} = \frac{p_\phi}{ml^2}$

$$H = \frac{p_\phi^2}{2ml^2} - mgl \cos(\phi)$$

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# Quantum Mechanics (in 20min)

Arguably, **physics** is ultimately about explaining experiments and forecasting measurement results.

Observables: Entities which are (actually) measured when an experiment is conducted on a system.

State: Entities which completely describe (or model) the system we are interested in.

Measurement establishes a relation between states and observables of a given system. Dynamics describes how observables and/or the state changes over time.

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# Postulates for Quantum Mechanics (ca 1950)

- The quantum **state** of a (free) particle is described by a (normalised) complex valued function:

$$\vec{\psi} \in L^2(x) \text{ i.e. } \int |\psi(x)|^2 dx = 1$$

- Two quantum states can be **superimposed**, i.e.

$$\alpha_1 \vec{\psi}_1 + \alpha_2 \vec{\psi}_2$$

- Any **observable**  $A$  is represented by a linear, self-adjoint operator  $\mathbf{A}$  on  $L^2(x)$ .
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$$\mathbf{A}\vec{\phi}_i = \lambda_i \vec{\phi}_i$$

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**Observables** and **states** of a system are represented by *hermitian* (i.e. self-adjoint) elements  $a$  of a  $C^*$ -algebra  $\mathcal{A}$  and by *states*  $w$  (i.e. normalised linear functionals) over this algebra.

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Quantum physics is often/sometimes counter-intuitive.

However, the standard mathematical model of (closed) quantum systems is relatively simple and just requires some basic **(complex) linear algebra**.

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# Quantum States and Notation

The state of a QM system is usually denoted by  $|x\rangle \in \mathcal{H}$ . The **inner product**  $\langle x|y\rangle$  of two vectors in  $\mathcal{H}$  – which is describing the *angle* between them – is very important in QM.

**P.A.M. Dirac** “invented” the Bra-Ket Notation based on the following simple facts:

Typewriters had no sub-scripts  $\vec{x}_i$   
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Simply “take inner product appart” to denote vectors in  $\mathcal{H}$ :

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# Conventions

## Physical Convention:

$$\langle x | \alpha y \rangle = \alpha \langle x | y \rangle$$

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Linear in first or second argument? In mathematics we have:

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- The **dynamics** of a closed system is described by the Schrödinger Equation:

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for the (self-adjoint) Hamiltonian  $\mathbf{H}$ .

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## Theorem

*For any self-adjoint operator  $\mathbf{A}$  the operator*

$$\exp(i\mathbf{A}) = e^{i\mathbf{A}} = \sum_{n=0}^{\infty} \frac{(i\mathbf{A})^n}{n!}$$

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# Adjoint Operator

For a matrix  $\mathbf{A} = (\mathbf{A}_{ij})$  its **transpose** matrix  $\mathbf{A}^T$  is defined as

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the **conjugate** matrix  $\mathbf{A}^*$  is defined by

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A square matrix/operator  $\mathbf{U}$  is called **unitary** if

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That means  $\mathbf{U}$ 's inverse is  $\mathbf{U}^\dagger = \mathbf{U}^{-1}$ . It also implies that  $\mathbf{U}$  is **invertible** and the inverse is easy to compute.

The postulates of **Quantum Mechanics** require that the **time evolution** to a quantum state, e.g. a qubit, are implemented via a unitary operator (as long as there is no measurement).

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# Unitary Operators

A square matrix/operator  $\mathbf{U}$  is called **unitary** if

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# Projections

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An operator  $\mathbf{P}$  on  $\mathbb{C}^n$  is called **projection** (or **idempotent**) iff

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An operator  $\mathbf{P}$  on  $\mathbb{C}^n$  is called **(orthogonal) projection** iff

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# Spectral Theorem

In the bra-ket notation we can represent a projection onto the sub-space generated by  $|x\rangle$  by the outer product  $\mathbf{P}_x = |x\rangle\langle x|$ .

## Theorem

*A self-adjoint operator  $\mathbf{A}$  (on a finite dimensional Hilbert space, e.g.  $\mathbb{C}^n$ ) can be represented uniquely as a linear combination*

$$\mathbf{A} = \sum_i \lambda_i \mathbf{P}_i$$

*with  $\lambda_i \in \mathbb{R}$  and  $\mathbf{P}_i$  the (orthogonal) projection onto the eigen-space generated by the eigen-vector  $|i\rangle$ , i.e.  $\mathbf{P}_i = |i\rangle\langle i|$*

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$$|x\rangle = \sum_i \mathbf{P}_i |x\rangle = \sum_i |i\rangle\langle i|x\rangle = \sum_i \langle i|x\rangle |i\rangle = \sum_i \alpha_i |i\rangle$$

With probability  $|\alpha_i|^2 = |\langle i|x\rangle|^2$  two things happen

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It is possible that for an eigen-value  $\lambda_i$  in the equation

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# Heisenberg and Schrödinger Picture

Describe the dynamics in terms of observables or states.

In particular if we consider not just pure (isolated) states, i.e. vectors in a Hilbert space, but instead probabilistic states which are represented by **density matrices**.

A density matrix  $\rho \in \mathcal{B}(\mathcal{H})$  is a Hermitian semi-positive definite matrix or operator with  $\text{trace}(\rho) = 1$ . Note that a given pure state  $|\psi\rangle$  can also be represented with density matrix  $|\psi\rangle\langle\psi|$ .

The quantum dynamics can be described as for **A observable** and  $\rho$  **state** (as density matrix/operator).

Schrödinger Picture:  $\varrho_t = \mathbf{U}(t)\varrho_0\mathbf{U}^*(t)$ .

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## Quantum Dynamics

For unitary transformations describing qubit dynamics:

$$\mathbf{U}^\dagger = \mathbf{U}^{-1}$$

The quantum dynamics is **invertible** or **reversible**

## Quantum Measurement

For projection operators involved in quantum measurement:

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# Dynamics of Programs

# pWhile – Syntax I

Full programs contain optional variable declarations:

$$\begin{array}{lcl} P & ::= & \mathbf{begin} \ S \ \mathbf{end} \\ & | & \mathbf{var} \ D \ \mathbf{begin} \ S \ \mathbf{end} \end{array}$$

Declarations are of the form:

$$\begin{array}{lcl} r & ::= & \mathbf{bool} \\ & | & \mathbf{int} \\ & | & \{ c_1, \dots, c_n \} \\ & | & \{ c_1 .. c_n \} \\ D & ::= & v : r \\ & | & v : r ; D \end{array}$$

with  $c_i$  (integer) constants and  $r$  denoting ranges.

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# pWhile – Syntax II

The syntax of statements  $S$  is as follows:

$S$	$::=$	<b>stop</b>
		<b>skip</b>
		$v := a$
		$v \text{ ?} = r$
		$S_1 ; S_2$
		<b>choose</b> $p_1 : S_1$ <b>or</b> $p_2 : S_2$ <b>ro</b>
		<b>if</b> $b$ <b>then</b> $S_1$ <b>else</b> $S_2$ <b>fi</b>
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# Evaluation of Expressions

$$\sigma \ni \mathbf{State} = \mathbf{Var} \rightarrow \mathbf{Z} \uplus \mathbf{T}$$

To illustrate approach consider only finite sub-range of  $\mathbf{Z}$ .

Evaluation  $\mathcal{E}$  of expressions  $e$  in state  $\sigma$ :

$$\mathcal{E}(n)\sigma = n$$

$$\mathcal{E}(v)\sigma = \sigma(v)$$

$$\mathcal{E}(a_1 \odot a_2)\sigma = \mathcal{E}(a_1)\sigma \odot \mathcal{E}(a_2)\sigma$$

$$\mathcal{E}(\mathbf{true})\sigma = \mathbf{tt}$$

$$\mathcal{E}(\mathbf{false})\sigma = \mathbf{ff}$$

$$\mathcal{E}(\mathbf{not } b)\sigma = \neg \mathcal{E}(b)\sigma$$

$$\dots = \dots$$

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# pWhile – SOS Semantics I

$$\mathbf{R0} \quad \langle \mathbf{skip}, \sigma \rangle \Rightarrow_1 \langle \mathbf{stop}, \sigma \rangle$$

$$\mathbf{R1} \quad \langle \mathbf{stop}, \sigma \rangle \Rightarrow_1 \langle \mathbf{stop}, \sigma \rangle$$

$$\mathbf{R2} \quad \langle v := e, \sigma \rangle \Rightarrow_1 \langle \mathbf{stop}, \sigma[v \mapsto \mathcal{E}(e)\sigma] \rangle$$

$$\mathbf{R3} \quad \langle v ?= r, \sigma \rangle \Rightarrow_{\frac{1}{|r|}} \langle \mathbf{stop}, \sigma[v \mapsto r_i \in r] \rangle$$

$$\mathbf{R4}_1 \quad \frac{\langle S_1, \sigma \rangle \Rightarrow_p \langle S'_1, \sigma' \rangle}{\langle S_1; S_2, \sigma \rangle \Rightarrow_p \langle S'_1; S_2, \sigma' \rangle}$$

$$\mathbf{R4}_2 \quad \frac{\langle S_1, \sigma \rangle \Rightarrow_p \langle \mathbf{stop}, \sigma' \rangle}{\langle S_1; S_2, \sigma \rangle \Rightarrow_p \langle S_2, \sigma' \rangle}$$

# pWhile – SOS Semantics II

**R5<sub>1</sub>**    $\langle \text{choose } p_1 : S_1 \text{ or } p_2 : S_2, \sigma \rangle \Rightarrow_{p_1} \langle S_1, \sigma \rangle$

**R5<sub>2</sub>**    $\langle \text{choose } p_1 : S_1 \text{ or } p_2 : S_2, \sigma \rangle \Rightarrow_{p_2} \langle S_2, \sigma \rangle$

**R6<sub>1</sub>**    $\langle \text{if } b \text{ then } S_1 \text{ else } S_2, \sigma \rangle \Rightarrow_1 \langle S_1, \sigma \rangle$       if  $\mathcal{E}(b)\sigma = \mathbf{tt}$

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**R7<sub>1</sub>**    $\langle \text{while } b \text{ do } S, \sigma \rangle \Rightarrow_1 \langle S; \text{ while } b \text{ do } S, \sigma \rangle$       if  $\mathcal{E}(b)\sigma = \mathbf{tt}$

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# Factorial

```
var
  m : {0..2};
  n : {0..2};

begin
  m := 1;
  while (n>1) do
    m := m*n;
    n := n-1;
  od;
  stop; # looping
end
```

# Multi Variable State

The problem we first consider is how to describe distributions over the cartesian product in order to represent the probabilities that two or more variables have certain values.

As we have  $\mathcal{D}(S) \subseteq \mathcal{V}(S)$  we investigate  $\mathcal{V}(S \times S)$ . In order to understand the relation between  $\mathcal{V}(S)$  and  $\mathcal{V}(S \times S)$  and in general  $\mathcal{V}(S^n)$  we need to consider the **tensor product**.

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# Multi Variable State

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As we have  $\mathcal{D}(S) \subseteq \mathcal{V}(S)$  we investigate  $\mathcal{V}(S \times S)$ . In order to understand the relation between  $\mathcal{V}(S)$  and  $\mathcal{V}(S \times S)$  and in general  $\mathcal{V}(S^n)$  we need to consider the **tensor product**.

Essential for the further treatment is the fact (more later) that

$$\mathcal{V}(S \times S) = \mathcal{V}(S) \otimes \mathcal{V}(S)$$

# Kronecker Product

Given a  $n \times m$  **matrix** **A** and a  $k \times l$  matrix **B**:

$$\mathbf{A} = \begin{pmatrix} a_{11} & \dots & a_{1m} \\ \vdots & \ddots & \vdots \\ a_{n1} & \dots & a_{nm} \end{pmatrix} \quad \mathbf{B} = \begin{pmatrix} b_{11} & \dots & b_{1l} \\ \vdots & \ddots & \vdots \\ b_{k1} & \dots & b_{kl} \end{pmatrix}$$

The **tensor** or **Kronecker product**  $\mathbf{A} \otimes \mathbf{B}$  is a  $nk \times ml$  matrix:

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Special cases are **square matrices** ( $n = m$  and  $k = l$ ) and **vectors** (row  $n = k = 1$ , column  $m = l = 1$ ).

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# Tensor Product Properties

The tensor product of  $n$  linear operators  $\mathbf{A}_1, \mathbf{A}_2, \dots, \mathbf{A}_n$  is associative (but in general not commutative) and has e.g. the following properties:

- 1  $(\mathbf{A}_1 \otimes \dots \otimes \mathbf{A}_n) \cdot (\mathbf{B}_1 \otimes \dots \otimes \mathbf{B}_n) = \mathbf{A}_1 \cdot \mathbf{B}_1 \otimes \dots \otimes \mathbf{A}_n \cdot \mathbf{B}_n$
- 2  $\mathbf{A}_1 \otimes \dots \otimes (\alpha \mathbf{A}_i) \otimes \dots \otimes \mathbf{A}_n = \alpha (\mathbf{A}_1 \otimes \dots \otimes \mathbf{A}_i \otimes \dots \otimes \mathbf{A}_n)$
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# Tensor Product Base

Every vector space has an algebraic base  $\{\mathbf{e}_i\}$

$$\mathbf{x} = x_1 \mathbf{e}_1 + x_2 \mathbf{e}_2 + \dots$$

This allows to specify vectors via coordinates  $\mathbf{x} = (x_1, x_2, \dots)$ .  
Base vectors are in this context simply of the form

$$\mathbf{e}_i = (e_{i1}, e_{i2}, \dots) \quad \text{with } e_{ij} = \begin{cases} 1 & \text{for } i = j \\ 0 & \text{otherwise} \end{cases}$$

The tensor product space  $\mathcal{V} \otimes \mathcal{W}$  can be seen as generated by (formal) tensors of the form  $\mathbf{v}_i \otimes \mathbf{w}_j$  with in  $\mathbf{v}_i \in \mathcal{V}$  and  $\mathbf{w}_j \in \mathcal{W}$  base vectors.

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# Classical and Probabilistic State

We have (always) a finite number  $v$  of **variables**.

**Classical state**  $\sigma \in \mathbf{State}$  given by:

$$\sigma \in \mathbf{State} = (\mathbf{Var} \rightarrow \mathbf{Value}) = \mathbf{Value}^v$$

For each variable we assume also a finite range of **values**.

**Probabilistic state**  $\mathbf{d}$  of a single variable is a distribution over possible values of the variable.

$$\mathbf{d} \in \mathcal{V}(\mathbf{Value}) = \{ (x_c)_{c \in \mathbf{Value}} \mid x_i \in \mathbb{R} \}$$

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# States and Tensor Products

For finite ranges we can represent distributions over cartesian product as an element in the tensor product in  $\mathcal{V}(\mathbf{Value})^{\otimes v}$ .

Probabilistic state  $\mathbf{d}$  of a all variables together

$$\begin{aligned}\mathbf{d} &\in \mathcal{V}(\mathbf{Var} \rightarrow \mathbf{Value}) = \\ &= \mathcal{V}(\mathbf{Value}_1 \times \mathbf{Value}_2 \times \dots \times \mathbf{Value}_v) = \\ &= \mathcal{V}(\mathbf{Value}_1) \otimes \mathcal{V}(\mathbf{Value}_2) \otimes \dots \otimes \mathcal{V}(\mathbf{Value}_v)\end{aligned}$$

For infinite value ranges we would need to consider measures. Product measures exist, for example, by Fubini's Theorem.

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# Probabilistic Control Flow

Consider the following (labelled) program:

```
1: while [ $z < 100$ ]1 do  
2:   choose2  $\frac{1}{3} : [x := 3]$ 3 or  $\frac{2}{3} : [x := 1]$ 4 ro  
3: od  
4: [stop]5
```

Its probabilistic control flow is given by:

$$\text{flow}(P) = \{\langle 1, 1, \underline{2} \rangle, \langle 1, 1, 5 \rangle, \langle 2, \frac{1}{3}, 3 \rangle, \langle 2, \frac{2}{3}, 4 \rangle, \langle 3, 1, 1 \rangle, \langle 4, 1, 1 \rangle\}.$$

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# Init — First Statement

$$\mathit{init}([\mathbf{skip}]^\ell) = \ell$$

$$\mathit{init}([\mathbf{stop}]^\ell) = \ell$$

$$\mathit{init}([\mathbf{v} := \mathbf{e}]^\ell) = \ell$$

$$\mathit{init}([\mathbf{v} ? = \mathbf{e}]^\ell) = \ell$$

$$\mathit{init}(S_1; S_2) = \mathit{init}(S_1)$$

$$\mathit{init}(\mathbf{choose}^\ell p_1 : S_1 \text{ or } p_2 : S_2) = \ell$$

$$\mathit{init}(\mathbf{if} [b]^\ell \text{ then } S_1 \text{ else } S_2) = \ell$$

$$\mathit{init}(\mathbf{while} [b]^\ell \text{ do } S) = \ell$$

# Final — Last Statements

$$final([\mathbf{skip}]^\ell) = \{\ell\}$$

$$final([\mathbf{stop}]^\ell) = \{\ell\}$$

$$final([\mathbf{v} := e]^\ell) = \{\ell\}$$

$$final([\mathbf{v} ?= e]^\ell) = \{\ell\}$$

$$final(S_1; S_2) = final(S_2)$$

$$final(\mathbf{choose}^\ell p_1 : S_1 \text{ or } p_2 : S_2) = final(S_1) \cup final(S_2)$$

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# Flow I — Control Transfer

$$\text{flow}([\mathbf{skip}]^\ell) = \emptyset$$

$$\text{flow}([\mathbf{stop}]^\ell) = \{\langle \ell, 1, \ell \rangle\}$$

$$\text{flow}([\mathbf{v} := e]^\ell) = \emptyset$$

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$$\begin{aligned} \text{flow}(S_1; S_2) &= \text{flow}(S_1) \cup \text{flow}(S_2) \cup \\ &\cup \{(\ell, 1, \text{init}(S_2)) \mid \ell \in \text{final}(S_1)\} \end{aligned}$$

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# Collecting Semantics

The **collecting semantics** of a program  $P$  is given by:

$$\mathbf{T}(P) = \sum_{\langle i, p_{ij}, j \rangle \in \mathcal{F}(P)} p_{ij} \cdot \mathbf{T}(\ell_i, \ell_j)$$

i.e. as a linear operator on  $\mathcal{V}(\mathbf{Value})^{\otimes \nu} \otimes \mathcal{V}(\mathbf{Lab})$ .

Local effects  $\mathbf{T}(\ell_i, \ell_j)$ : Data Update  $\mathbf{N}$  + Control Step  $\mathbf{M}$

$$\mathbf{T}(\ell_i, \ell_j) = \mathbf{N}_i \otimes \mathbf{M}_{ij} = \mathbf{N}_{i1} \otimes \mathbf{N}_{i2} \otimes \dots \otimes \mathbf{N}_{i\nu} \otimes \mathbf{M}_{ij}$$

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i.e. as a linear operator on  $\mathcal{V}(\mathbf{Value})^{\otimes \nu} \otimes \mathcal{V}(\mathbf{Lab})$ .

Local effects  $\mathbf{T}(\ell_i, \ell_j)$ : Data Update  $\mathbf{N}$  + Control Step  $\mathbf{M}$

$$\mathbf{T}(\ell_i, \ell_j) = \mathbf{N}_i \otimes \mathbf{M}_{ij} = \mathbf{N}_{i1} \otimes \mathbf{N}_{i2} \otimes \dots \otimes \mathbf{N}_{i\nu} \otimes \mathbf{M}_{ij}$$



# Local Transfer Operators

$$\mathbf{T}(\ell_1, \ell_2) = \mathbf{I} \otimes \mathbf{E}(\ell_1, \ell_2) \quad \text{for } [\mathbf{skip}]^{\ell_1}$$

$$\mathbf{T}(\ell, \ell) = \mathbf{I} \otimes \mathbf{E}(\ell, \ell) \quad \text{for } [\mathbf{stop}]^{\ell}$$

$$\mathbf{T}(\ell_1, \ell_2) = \mathbf{U}(v \leftarrow e) \otimes \mathbf{E}(\ell_1, \ell_2) \quad \text{for } [v := e]^{\ell_1}$$

$$\mathbf{T}(\ell_1, \ell_2) = \left( \frac{1}{|r|} \sum_{c \in r} \mathbf{U}(v := c) \right) \otimes \mathbf{E}(\ell_1, \ell_2) \quad \text{for } [v ?= r]^{\ell_1}$$

$$\mathbf{T}(\ell, \ell_k) = \mathbf{I} \otimes \mathbf{E}(\ell, \ell_k) \quad \text{for } [\mathbf{choose}]^{\ell}$$

$$\mathbf{T}(\ell, \underline{\ell}_t) = \mathbf{P}(b = \mathbf{tt}) \otimes \mathbf{E}(\ell, \ell_t) \quad \text{for } [b]^{\ell}$$

$$\mathbf{T}(\ell, \ell_f) = \mathbf{P}(b = \mathbf{ff}) \otimes \mathbf{E}(\ell, \ell_f) \quad \text{for } [b]^{\ell}$$

# Trivial Operators

Matrix Units – Represent a single transition

$$(\mathbf{E}(m, n))_{ij} = \begin{cases} 1 & \text{if } m = i \wedge n = j \\ 0 & \text{otherwise.} \end{cases}$$

Identity – Represents “no change” transition

$$(\mathbf{I})_{ij} = \begin{cases} 1 & \text{if } i = j \\ 0 & \text{otherwise.} \end{cases}$$

# Tests Operators and Filters

Select a certain value  $c \in \mathbf{Value}$ :

$$(\mathbf{P}(c))_{ij} = \begin{cases} 1 & \text{if } i = c = j \\ 0 & \text{otherwise.} \end{cases}$$

Select a certain classical state  $\sigma \in \mathbf{State}$ :

$$\mathbf{P}(\sigma) = \bigotimes_{i=1}^V \mathbf{P}(\sigma(v_i))$$

Select states where expression  $e = a \mid b$  evaluates to  $c$ :

$$\mathbf{P}(e = c) = \sum_{\mathcal{E}(e)\sigma=c} \mathbf{P}(\sigma)$$

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# Selection via Projections

Filtering out *relevant* configurations, i.e. only those which fulfill a certain condition. Use diagonal matrix  $\mathbf{P}$ :

$$(\mathbf{P})_{ii} = \begin{cases} 1 & \text{if condition holds for } C_i \\ 0 & \text{otherwise.} \end{cases}$$

$$\begin{pmatrix} C_1 \\ C_2 \\ C_3 \\ C_4 \\ C_5 \\ C_6 \end{pmatrix}^t \cdot \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 \\ C_2 \\ C_3 \\ 0 \\ C_5 \\ 0 \end{pmatrix}^t$$

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# Update Operators

For all initial values change to constant  $c \in \mathbf{Value}$ :

$$(\mathbf{U}(c))_{ij} = \begin{cases} 1 & \text{if } j = c \\ 0 & \text{otherwise.} \end{cases}$$

Set value of the  $k$ th variable  $v_k \in \mathbf{Var}$  to constant  $c \in \mathbf{Value}$ :

$$\mathbf{U}(v_k \leftarrow c) = \left( \bigotimes_{i=1}^{k-1} \mathbf{I} \right) \otimes \mathbf{U}(c) \otimes \left( \bigotimes_{i=k+1}^v \mathbf{I} \right)$$

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# Program Approximation

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A general approach towards problems and attempts to solve them could be described as follows:

- If the problem is to **difficult**
  - formulate a **simplified** version,
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# Probabilistic Program Analysis

Possible aims of **Static Program Analysis**:

- **Safe Approximations**:  
Correct under all circumstances.
- **Good/Close Estimates**:  
Fix it (at runtime) if there is a problem.

With modern computer architectures some **compile time** tasks (type checking, threading, etc.) become **runtime** features.

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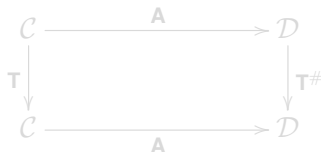
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# Semantical Abstraction

Consider a **Concrete Domain**  $\mathcal{C}$  and an **Abstract Domain**  $\mathcal{D}$ :

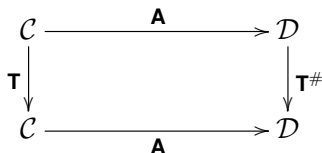


With an **abstraction**  $A : \mathbf{C} \rightarrow \mathbf{D}$  and a **concretisation**  $G : \mathbf{D} \rightarrow \mathbf{C}$ :

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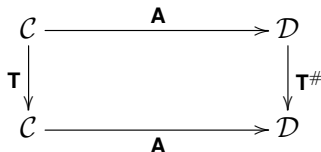
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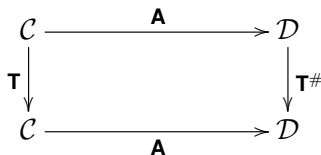


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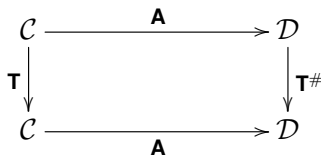
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Abstract Interpretation:  $(\mathbf{A}, \mathbf{G})$  form a **Galois Connection**.

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Probabilistic Abst.Int.:  $(\mathbf{A}, \mathbf{G})$  **Moore-Penrose Pseudo-Inverse**.

# Galois Connections

## Definition

Let  $\mathcal{C} = (\mathcal{C}, \leq)$  and  $\mathcal{D} = (\mathcal{D}, \sqsubseteq)$  be two partially ordered set. If there are two functions  $\alpha : \mathcal{C} \rightarrow \mathcal{D}$  and  $\gamma : \mathcal{D} \rightarrow \mathcal{C}$  such that for all  $c \in \mathcal{C}$  and all  $d \in \mathcal{D}$ :

$$c \leq \gamma(d) \text{ iff } \alpha(c) \sqsubseteq d,$$

then  $(\mathcal{C}, \alpha, \gamma, \mathcal{D})$  form a **Galois connection**.

# Moore-Penrose Pseudo-Inverse I

## Definition

Let  $\mathcal{C}$  and  $\mathcal{D}$  be two Hilbert spaces and  $\mathbf{A} : \mathcal{C} \rightarrow \mathcal{D}$  a bounded linear map. A bounded linear map  $\mathbf{A}^\dagger = \mathbf{G} : \mathcal{D} \rightarrow \mathcal{C}$  is the **Moore-Penrose pseudo-inverse** of  $\mathbf{A}$  iff

$$(i) \quad \mathbf{A} \circ \mathbf{G} = \mathbf{P}_A,$$

$$(ii) \quad \mathbf{G} \circ \mathbf{A} = \mathbf{P}_G,$$

where  $\mathbf{P}_A$  and  $\mathbf{P}_G$  denote orthogonal projections onto the ranges of  $\mathbf{A}$  and  $\mathbf{G}$ .

# (Orthogonal) Projections – Idempotents

On finite dimensional vector (Hilbert) spaces we have an **inner product**  $\langle \cdot, \cdot \rangle$ . This allows us to define an **adjoint** via:

$$\langle \mathbf{A}(x), y \rangle = \langle x, \mathbf{A}^*(y) \rangle$$

- An operator  $\mathbf{A}$  is **self-adjoint** if  $\mathbf{A} = \mathbf{A}^*$ .
- An operator  $\mathbf{A}$  is **positive**, i.e.  $\mathbf{A} \supseteq 0$ , if there exists an operator  $\mathbf{B}$  such that  $\mathbf{A} = \mathbf{B}^*\mathbf{B}$ .
- An **(orthogonal) projection** is a self-adjoint  $\mathbf{E}$  with  $\mathbf{E}\mathbf{E} = \mathbf{E}$ .

Projections identify (closed) sub-spaces  $Y_{\mathbf{E}} = \{\mathbf{E}x \mid x \in \mathcal{V}\}$ .

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# Moore-Penrose Pseudo-Inverse II

## Definition

An operator  $\mathbf{A} \in \mathcal{B}(\mathcal{H})$  is **Moore-Penrose invertible** if there exists an element  $\mathbf{G} \in \mathcal{B}(\mathcal{H})$  such that:

- (i)  $\mathbf{AGA} = \mathbf{A}$ ,
- (ii)  $\mathbf{GAG} = \mathbf{G}$ ,
- (iii)  $(\mathbf{AG})^* = \mathbf{AG}$ ,
- (iv)  $(\mathbf{GA})^* = \mathbf{GA}$ .

If it exists  $\mathbf{G} = \mathbf{A}^\dagger$  is called **Moore-Penrose pseudo-inverse**.

# Galois Connection II

## Definition

Let  $\mathcal{C} = (\mathcal{C}, \leq_{\mathcal{C}})$  and  $\mathcal{D} = (\mathcal{D}, \leq_{\mathcal{D}})$  be two partially ordered sets with two order-preserving functions  $\alpha : \mathcal{C} \mapsto \mathcal{D}$  and  $\gamma : \mathcal{D} \mapsto \mathcal{C}$ . Then  $(\mathcal{C}, \alpha, \gamma, \mathcal{D})$  form a **Galois connection** iff

- (i)  $\alpha \circ \gamma$  is **reductive** i.e.  $\forall d \in \mathcal{D}, \alpha \circ \gamma(d) \leq_{\mathcal{D}} d$ ,
- (ii)  $\gamma \circ \alpha$  is **extensive** i.e.  $\forall c \in \mathcal{C}, c \leq_{\mathcal{C}} \gamma \circ \alpha(c)$ .

## Proposition

*Let  $(\mathcal{C}, \alpha, \gamma, \mathcal{D})$  be a Galois connection. Then  $\alpha$  and  $\gamma$  are **quasi-inverse**, i.e.*

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# Examples of Abstractions

**Parity Abstraction** operator on  $\mathcal{V}(\{1, \dots, n\})$  (with  $n$  even):

$$\mathbf{A}_p = \begin{pmatrix} 1 & 0 \\ 0 & 1 \\ 1 & 0 \\ 0 & 1 \\ \vdots & \vdots \\ 0 & 1 \end{pmatrix} \quad \mathbf{A}_p^\dagger = \begin{pmatrix} \frac{2}{n} & 0 & \frac{2}{n} & 0 & \dots & 0 \\ 0 & \frac{2}{n} & 0 & \frac{2}{n} & \dots & \frac{2}{n} \end{pmatrix}$$

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**Sign Abstraction** operator on  $\mathcal{V}(\{-n, \dots, 0, \dots, n\})$ :

$$\mathbf{A}_S = \begin{pmatrix} 1 & 0 & 0 \\ \vdots & \vdots & \vdots \\ 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \\ \vdots & \vdots & \vdots \\ 0 & 0 & 1 \end{pmatrix} \quad \mathbf{A}_S^\dagger = \begin{pmatrix} \frac{1}{n} & \dots & \frac{1}{n} & 0 & 0 & \dots & 0 \\ 0 & \dots & 0 & 1 & 0 & \dots & 0 \\ 0 & \dots & 0 & 0 & \frac{1}{n} & \dots & \frac{1}{n} \end{pmatrix}$$



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# Lifting of an extraction function $\alpha : \mathcal{C} \mapsto \mathcal{D}$

**Power Set lifting** to an abstraction function  $\tilde{\alpha} : \mathcal{P}(\mathcal{C}) \rightarrow \mathcal{P}(\mathcal{D})$

$$\tilde{\alpha}(\{c_1, c_2, \dots\}) = \{\alpha(c_1), \alpha(c_2), \dots\}$$

**Vector Space lifting** to an abstraction function  $\vec{\alpha} : \mathcal{V}(\mathcal{C}) \rightarrow \mathcal{V}(\mathcal{D})$

$$\vec{\alpha}(p_1 \cdot \vec{c}_1 + p_2 \cdot \vec{c}_2 + \dots) = p_1 \cdot \alpha(c_1) + p_2 \cdot \alpha(c_2) \dots$$

**Support Set:**  $\text{supp} : \mathcal{V}(\mathcal{C}) \rightarrow \mathcal{P}(\mathcal{C})$

$$\text{supp}(\vec{x}) = \{c_i \mid \langle c_i, p_i \rangle \in \vec{x} \text{ and } p_i \neq 0\}$$

**Uniform Distribution:**  $\text{vec} : \mathcal{P}(\mathcal{C}) \rightarrow \mathcal{V}(\mathcal{C})$

$$\text{vec}(\tilde{x}) = \{\langle c_i, 1/|\tilde{x}| \rangle\}$$

# Lifting of an extraction function $\alpha : \mathcal{C} \mapsto \mathcal{D}$

**Power Set lifting** to an abstraction function  $\tilde{\alpha} : \mathcal{P}(\mathcal{C}) \rightarrow \mathcal{P}(\mathcal{D})$

$$\tilde{\alpha}(\{c_1, c_2, \dots\}) = \{\alpha(c_1), \alpha(c_2), \dots\}$$

**Vector Space lifting** to an abstraction function  $\vec{\alpha} : \mathcal{V}(\mathcal{C}) \rightarrow \mathcal{V}(\mathcal{D})$

$$\vec{\alpha}(p_1 \cdot \vec{c}_1 + p_2 \cdot \vec{c}_2 + \dots) = p_1 \cdot \alpha(c_1) + p_2 \cdot \alpha(c_2) \dots$$

**Support Set:**  $\text{supp} : \mathcal{V}(\mathcal{C}) \rightarrow \mathcal{P}(\mathcal{C})$

$$\text{supp}(\vec{x}) = \{c_i \mid \langle c_i, p_i \rangle \in \vec{x} \text{ and } p_i \neq 0\}$$

**Uniform Distribution:**  $\text{vec} : \mathcal{P}(\mathcal{C}) \rightarrow \mathcal{V}(\mathcal{C})$

$$\text{vec}(\tilde{x}) = \{\langle c_i, 1/|\tilde{x}| \rangle\}$$

# Relation between Abstractions [PPDP00]

## Proposition

Let  $\vec{\alpha}$  be a *probabilistic abstraction* function and let  $\vec{\gamma}$  be its Moore-Penrose pseudo-inverse.

Then  $\vec{\gamma} \circ \vec{\alpha}$  is *extensive* with respect to the inclusion on the support sets of vectors in  $\mathcal{V}(\mathcal{C})$ , i.e.  $\forall \vec{x} \in \mathcal{V}(\mathcal{C})$ ,

$$\text{supp}(\vec{x}) \subseteq \text{supp}(\vec{\gamma} \circ \vec{\alpha}(\vec{x})).$$

# Least Square Approximation

Given a linear equation

$$x\mathbf{A} = b$$

it has either (i) a (unique) **solution**  $\bar{x}$ , or (ii) the **residual**

$$r_x = b - x\mathbf{A}$$

is non-zero for all  $x$ .

The (unique) **least-square solution**  $\bar{x}$ , i.e. for which the residual  $\|b - \bar{x}\mathbf{A}\|$  is minimal, can be obtained using the Moore-Penrose pseudo-inverse:

$$\bar{x} = b\mathbf{A}^\dagger$$

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# Abstract LOS Semantics

Moore-Penrose Pseudo-Inverse of a Tensor Product is simply

$$(\mathbf{A}_1 \otimes \mathbf{A}_2 \otimes \dots \otimes \mathbf{A}_n)^\dagger = \mathbf{A}_1^\dagger \otimes \mathbf{A}_2^\dagger \otimes \dots \otimes \mathbf{A}_n^\dagger$$

Via linearity we can construct  $\mathbf{T}^\#$  in the same way as  $\mathbf{T}$ , i.e

$$\mathbf{T}^\#(P) = \sum_{\langle i, p_{ij}, j \rangle \in \mathcal{F}(P)} p_{ij} \cdot \mathbf{T}^\#(\ell_i, \ell_j)$$

with local abstraction of individual variables:

$$\mathbf{T}^\#(\ell_i, \ell_j) = (\mathbf{A}_1^\dagger \mathbf{N}_{i1} \mathbf{A}_1) \otimes (\mathbf{A}_2^\dagger \mathbf{N}_{i2} \mathbf{A}_2) \otimes \dots \otimes (\mathbf{A}_v^\dagger \mathbf{N}_{iv} \mathbf{A}_v) \otimes \mathbf{M}_{ij}$$

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# Proof Argument

$$\begin{aligned}\mathbf{T}^\# &= \mathbf{A}^\dagger \mathbf{T} \mathbf{A} \\ &= \mathbf{A}^\dagger \left( \sum_{i,j} \mathbf{T}(i,j) \right) \mathbf{A} \\ &= \sum_{i,j} \mathbf{A}^\dagger \mathbf{T}(i,j) \mathbf{A} \\ &= \sum_{i,j} \left( \bigotimes_k \mathbf{A}_k \otimes \mathbf{I} \right)^\dagger \mathbf{T}(i,j) \left( \bigotimes_k \mathbf{A}_k \otimes \mathbf{I} \right) \\ &= \sum_{i,j} \left( \bigotimes_k \mathbf{A}_k \otimes \mathbf{I} \right)^\dagger \left( \bigotimes_k \mathbf{N}_{ik} \otimes \mathbf{M}_{ij} \right) \left( \bigotimes_k \mathbf{A}_k \otimes \mathbf{I} \right) \\ &= \sum_{i,j} \left( \bigotimes_k (\mathbf{A}_k^\dagger \mathbf{N}_{ik} \mathbf{A}_k) \otimes \mathbf{M}_{ij} \right)\end{aligned}$$

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# Example: Factorial

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1:  $[m \leftarrow 1]^1$ ;  
2: while  $[n > 1]^2$  do  
3:    $[m \leftarrow m \times n]^3$ ;  
4:    $[n \leftarrow n - 1]^4$   
5: od  
6:  $[\text{stop}]^5$ 
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$$\begin{aligned} \mathbf{T} &= \mathbf{U}(m \leftarrow 1) \otimes \mathbf{E}(1, 2) \\ &+ \mathbf{P}(n > 1) \otimes \mathbf{E}(2, 3) \\ &+ \mathbf{P}(n \leq 1) \otimes \mathbf{E}(2, 5) \\ &+ \mathbf{U}(m \leftarrow m \times n) \otimes \mathbf{E}(3, 4) \\ &+ \mathbf{U}(n \leftarrow n - 1) \otimes \mathbf{E}(4, 2) \\ &+ \mathbf{I} \otimes \mathbf{E}(5, 5) \end{aligned}$$

The abstract versions of the local filters and updates, e.g.  $\mathbf{P}^\#(n > 1)$ ,  $\mathbf{U}^\#(m \leftarrow m \times n)$ ,  $\mathbf{U}^\#(n \leftarrow n - 1)$  etc. justify our previous ad hoc analysis.



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# Abstract Semantics

**Abstraction:**  $\mathbf{A} = \mathbf{A}_p \otimes \mathbf{I}$ , i.e. *m abstract* (parity) but *n concrete*.

$$\begin{aligned} \mathbf{T}^\# &= \mathbf{U}^\#(m \leftarrow 1) \otimes \mathbf{E}(1, 2) \\ &+ \mathbf{P}^\#(n > 1) \otimes \mathbf{E}(2, 3) \\ &+ \mathbf{P}^\#(n \leq 1) \otimes \mathbf{E}(2, 5) \\ &+ \mathbf{U}^\#(m \leftarrow m \times n) \otimes \mathbf{E}(3, 4) \\ &+ \mathbf{U}^\#(n \leftarrow n - 1) \otimes \mathbf{E}(4, 2) \\ &+ \mathbf{I}^\# \otimes \mathbf{E}(5, 5) \end{aligned}$$

# Abstract Semantics

$$\begin{aligned} \mathbf{U}^\#(m \leftarrow i) = \\ = \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 & 0 & 0 & \dots & 0 \\ 0 & 1 & 0 & 0 & \dots & 0 \\ 0 & 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 0 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & & \dots & 1 \end{pmatrix} \end{aligned}$$

# Abstract Semantics

$$\begin{aligned} \mathbf{U}^\#(n \leftarrow n - 1) &= \\ &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 0 & 0 & 0 & 0 & \dots & 0 \\ 1 & 0 & 0 & 0 & \dots & 0 \\ 0 & 1 & 0 & 0 & \dots & 0 \\ 0 & 0 & 1 & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & 0 & \dots & 0 \end{pmatrix} \end{aligned}$$

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$$\mathbf{P}^\#(n > 1)) \\ = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 0 & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 0 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & 0 & \dots & 1 \end{pmatrix}$$

# Abstract Semantics

$$\begin{aligned} \mathbf{P}^\#(n \leq 1) &= \\ &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 & 0 & 0 & \dots & 0 \\ 0 & 1 & 0 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & 0 & \dots & 0 \end{pmatrix} \end{aligned}$$

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 & + \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & 0 & \dots & \ddots \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 0 & 0 & 0 & 0 & \dots & 0 \\ 0 & 1 & 0 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & 0 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & 0 & \dots & \ddots \end{pmatrix}
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# Implementation

Implementation of concrete and abstract semantics of **Factorial** using **octave**. **Ranges:**  $n \in \{1, 2, \text{max}\}$  and  $m \in \{1, 2, \text{max!}\}$ .

$n$	$\dim(\mathbf{T}(F))$	$\dim(\mathbf{T}^\#(F))$
2	45	30
3	140	40
4	625	50
5	3630	60
6	25235	70
7	201640	80
8	1814445	90
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Using **uniform** initial distributions  $\mathbf{d}_0$  for  $n$  and  $m$ .

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# Scaleability

The abstract probabilities for  $m$  being **even** or **odd** when we execute the abstract program for various maximal  $n$  values are:

$n$	<b>even</b>	<b>odd</b>
10	0.81818	0.18182
100	0.98019	0.019802
1000	0.99800	0.0019980
10000	0.99980	0.00019998

# The End

# Bibliography



Morris W. Hirsch, Stephen Smale, and Robert L. Devaney.  
*Differential Equations, Dynamical Systems and An  
Introduction to Chaos.*  
Elsevier, 2004.



L.D. Landau and E.M. Lifschitz.  
*Mechanik.*  
Akademie-Verlag, Berlin, 1981.



Barry Simon.  
*Representation of Finite and Compact Groups*, volume 10  
of *Graduate Studies in Mathematics.*  
AMS, 1996.